

Odd frequency pairing in superconducting heterostructures

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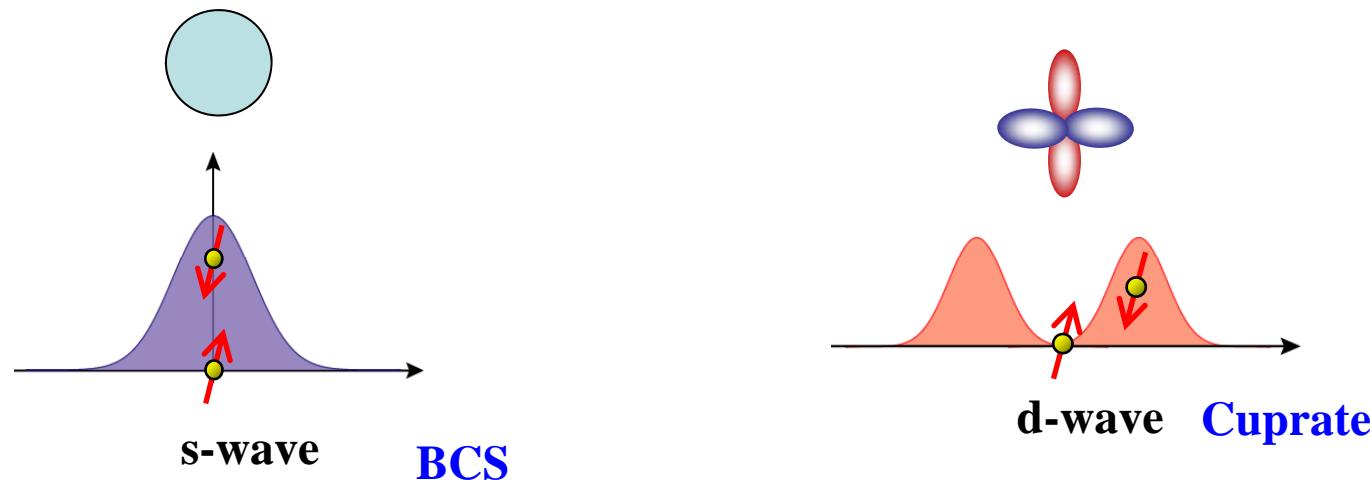
Akita University, Japan

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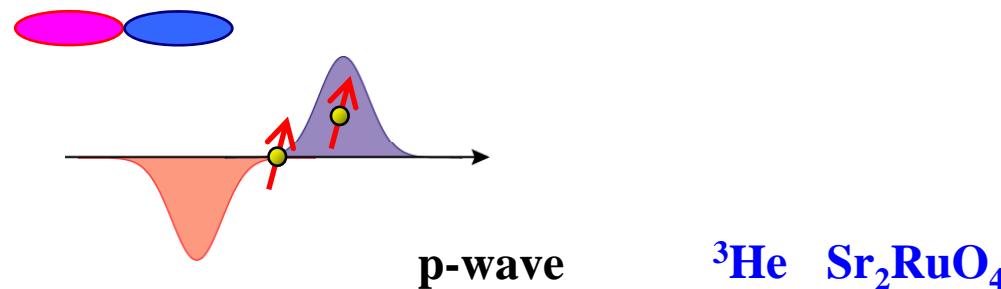
- (1) What is odd-frequency pairing
- (2) Normal metal / Superconductor junctions
- (3) Vortices in superconductors

Conventional Classification of Symmetry of Cooper pair

Spin-singlet Cooper pair \longrightarrow **Even Parity**



Spin-triplet Cooper pair \longrightarrow **Odd Parity**



Odd-frequency pairing

Fermi-Dirac statistics

Symmetry of pair wave functions:

$$\mathbf{k} \otimes \sigma \otimes \omega = \text{odd}$$

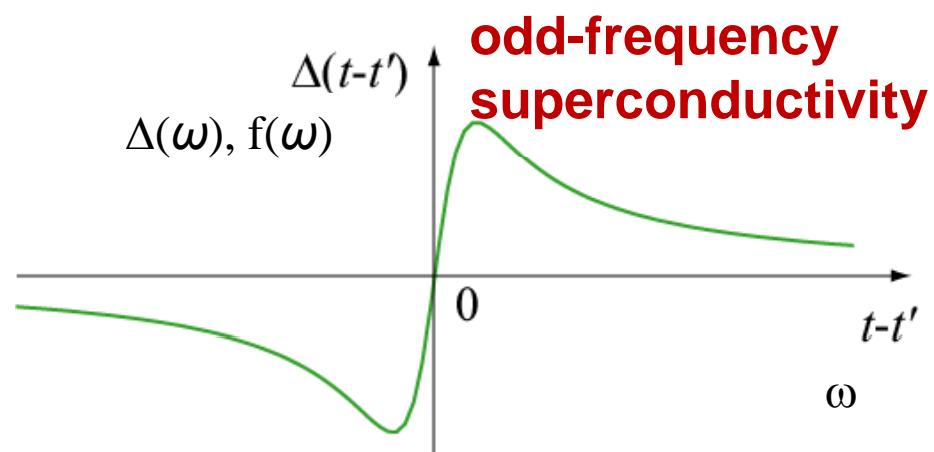
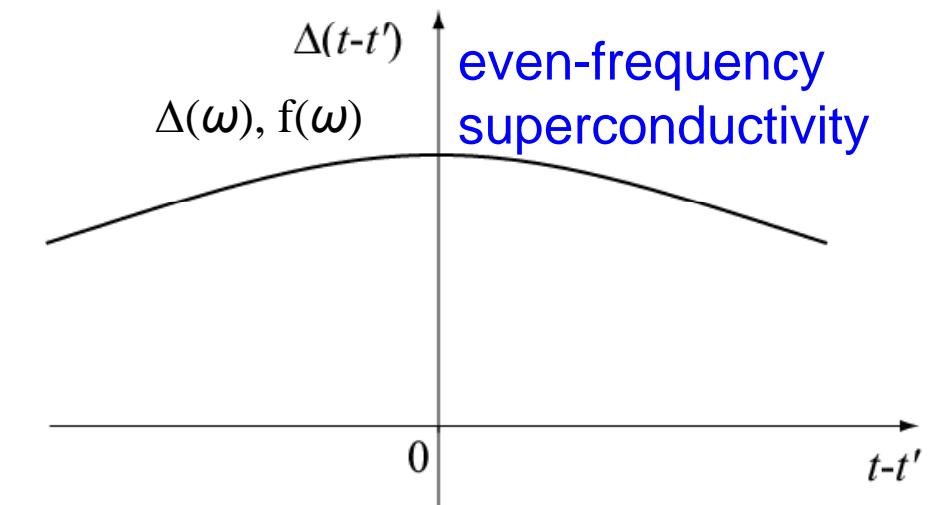
Momentum x Spin x Frequency



Berezinskii
(1974):

Spin-triplet s-wave

Balatsky&Abrahams
(1992):
Spin-singlet p-wave



ω , Matsubara frequency

Odd-frequency pairing state

(1) Odd-frequency pairing (**pair potential**, gap function) in uniform bulk system (odd-frequency superconductor)

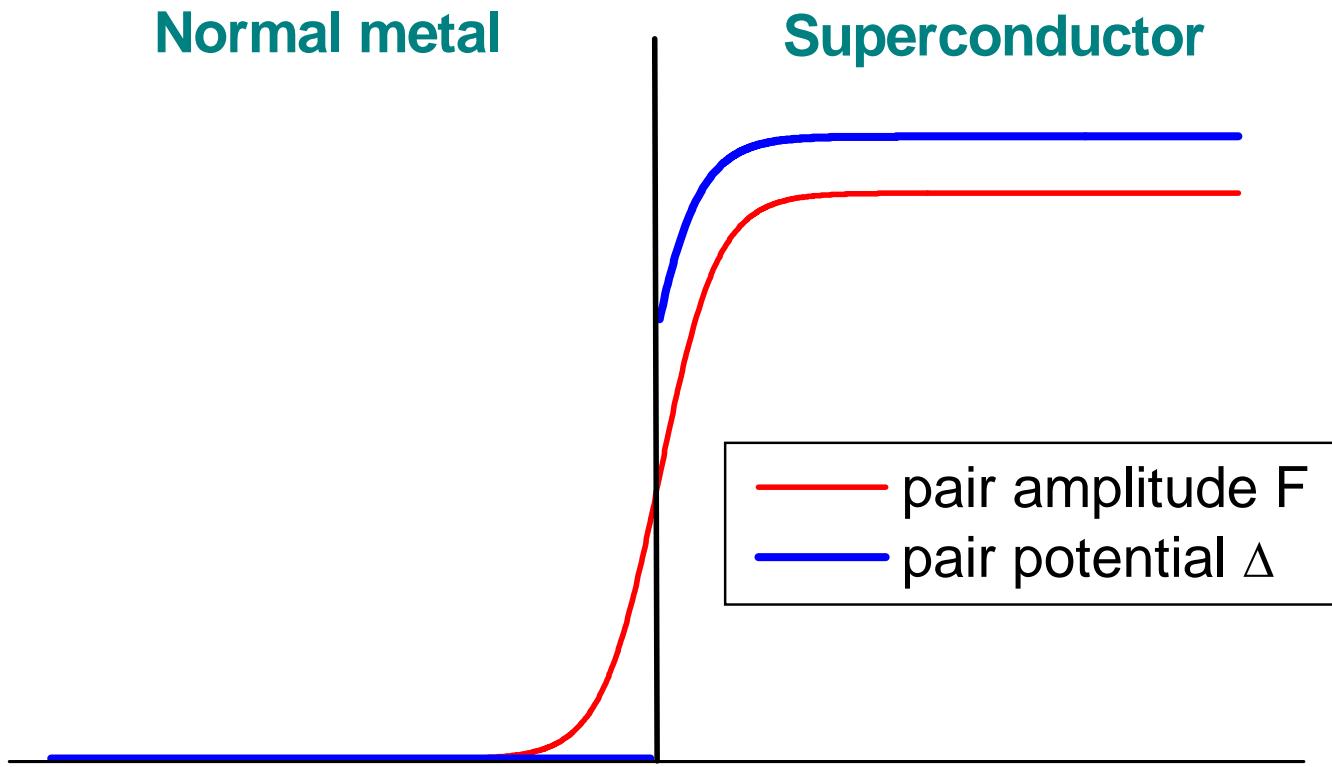
Uniform (bulk) system:

$$\Delta(\mathbf{k}, i\omega_n) = -T \sum_{k',m} V(\mathbf{k}-\mathbf{k}', i\omega_n - i\omega_m) F(\mathbf{k}', i\omega_m)$$



Pair potential **electron-interaction** **Pair amplitude**
Energy Gap

(2) Odd-frequency pairing state (**pair amplitude**) in superconducting junctions



Weak coupling BCS:
$$\Delta = \lambda T \sum_m F(i\omega_m)$$

Pair amplitude (pair correlation)

$$\begin{aligned} F_{\alpha,\beta}(\mathbf{r}_1 t_1, \mathbf{r}_2 t_2) &= -i \left\langle \mathcal{T} \psi_\alpha(\mathbf{r}_1 t_1) \psi_\beta(\mathbf{r}_2 t_2) \right\rangle \\ &= -i\theta(t_1-t_2) \left\langle \psi_\alpha(\mathbf{r}_1 t_1) \psi_\beta(\mathbf{r}_2 t_2) \right\rangle + i\theta(t_2-t_1) \left\langle \psi_\beta(\mathbf{r}_2 t_2) \psi_\alpha(\mathbf{r}_1 t_1) \right\rangle \end{aligned}$$

Exchange of two electrons

$$F_{\alpha,\beta}(\mathbf{r}_1 t_1, \mathbf{r}_2 t_2) = -F_{\beta,\alpha}(\mathbf{r}_2 t_2, \mathbf{r}_1 t_1)$$

Fermi-Dirac statistics

Pair amplitude

Exchange of time

Even-frequency pairing (conventional pairing)

$$F_{\alpha,\beta}(\mathbf{r}_1 t_1, \mathbf{r}_2 t_2) = F_{\alpha,\beta}(\mathbf{r}_1 t_2, \mathbf{r}_2 t_1)$$

Odd-frequency pairing

$$F_{\alpha,\beta}(\mathbf{r}_1 t_1, \mathbf{r}_2 t_2) = -F_{\alpha,\beta}(\mathbf{r}_1 t_2, \mathbf{r}_2 t_1)$$

Symmetry of the pair amplitude

+ symmetric, – anti-symmetric

	Frequency (time)	Spin	Orbital	Total	
ESE	+ (even)	– (singlet)	+ (even)	–	BCS
ETO	+ (even)	+ (triplet)	– (odd)	–	Cuprate
OTE	– (odd)	+ (triplet)	+ (even)	–	^3He
OSO	– (odd)	– (singlet)	– (odd)	–	Sr_2RuO_4

ESE (Even-frequency spin-singlet even-parity)

ETO (Even-frequency spin-triplet odd-parity)

OTE (Odd-frequency spin-triplet even-parity) Berezinskii

OSO (Odd-frequency spin-singlet odd-parity) Balatsky, Abrahams

Previous studies of odd-frequency pairing

Bulk state (Pair potential, Gap function)

Berezinskii (1974)

Balatsky, Abrahams, Schrieffer, Scalapino(1992-1993)

Zachar, Kivelson, Emery (1996)

Coleman, Miranda, Tsvelik (1997)

Vojta, Dagotto (1999)

Fuseya, Kohno, Miyake (2003)

Shigeta, Onari, Yada, Tanaka (2009)

Junction (No pair potential)

Induced odd-frequency pair amplitude in ferromagnet attached to spin-singlet s-wave superconductor

Bergeret, Efetov, Volkov, (2001)

- **Odd-frequency pairing state** is possible in inhomogeneous superconductors even for conventional even-frequency paring in the bulk

Broken spin rotation symmetry or spatial invariance symmetry can induce odd-frequency pairing state:

- **ferromagnet/superconducor junctions:**

Bergeret,Volkov&Efetov, 2001

- **non-uniform systems:**

Junctions: Tanaka&Golubov, 2007; Eschrig&Lofwander, 2007

Vortices: Yokoyama *et al.*, 2008; Tanuma *et al.*, 2009)

Contents

- (1) What is odd-frequency pairing
- (2) Ballistic normal metal junctions
- (3) Vortices in superconductors

Ballistic junction

**Ballistic
Normal metal
(semi-infinite)**

**Superconductor
(semi-infinite)**

Y. Tanaka, A. Golubov, S. Kashiwaya, and M. Ueda
Phys. Rev. Lett. 99 037005 (2007)

M. Eschrig, T. Lofwander, Th. Champel, J.C. Cuevas and G. Schon
J. Low Temp. Phys 147 457(2007)

Eilenberger equation

(explicitly denote direction of motion)

$$\mp iv_F x \partial_x f_{1\pm} = 2\omega_n f_{2\pm} - 2\bar{\Delta}_\pm(x)g_\pm$$

$$\mp iv_F x \partial_x g_\pm = 2\bar{\Delta}_\pm(x)f_{1\pm},$$

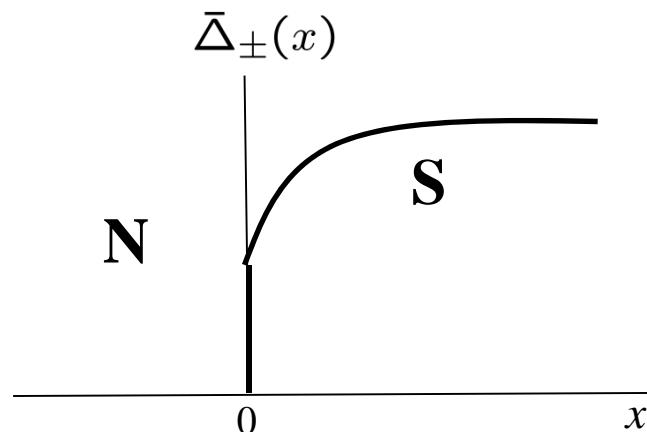
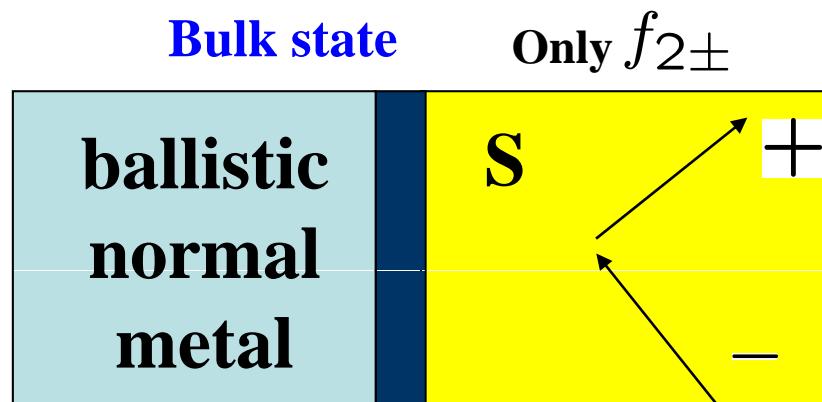
$$\mp iv_F x \partial_x f_{2\pm} = -2\omega_n f_{1\pm},$$

Pair potential	$\bar{\Delta}_\pm(x)$
Quasiparticle function	g_\pm
Pair amplitudes	$f_{1\pm}, f_{2\pm}$

$$f_{1\pm}^2 + f_{2\pm}^2 + g_\pm^2 = 1,$$

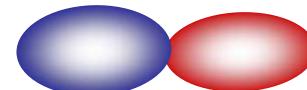
$$\bar{\Delta}_\pm(x) = \Delta(x)\Phi_\pm(\theta).$$

$\Phi_\pm(\theta)$ Form factor

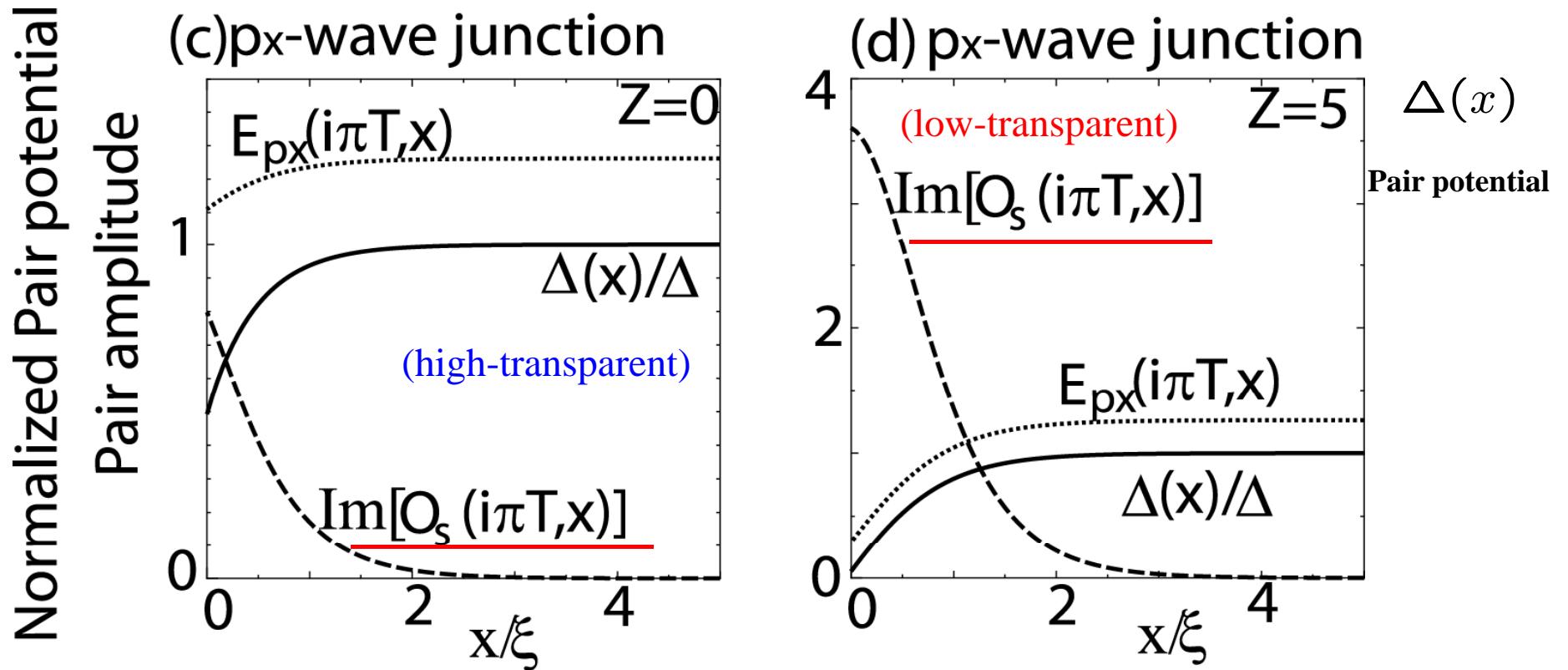


Normal metal

**spin-triplet p-wave
superconductor**



Symmetry of the bulk pair potential is ETO



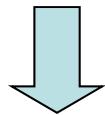
$E_{px}(i\omega_n, x)$ px-wave component of ETO pair amplitude
 $O_s(i\omega_n, x)$ s-wave component of OTE pair amplitude

ETO (Even-frequency spin-triplet odd-parity)

OTE (Odd-frequency spin-triplet even-parity) Y. Tanaka, et al PRL 99 037005 (2007)

Underlying physics

Near the interface, **even and odd-parity** pairing states (pair amplitude) can mix due to the **breakdown of the translational symmetry**.

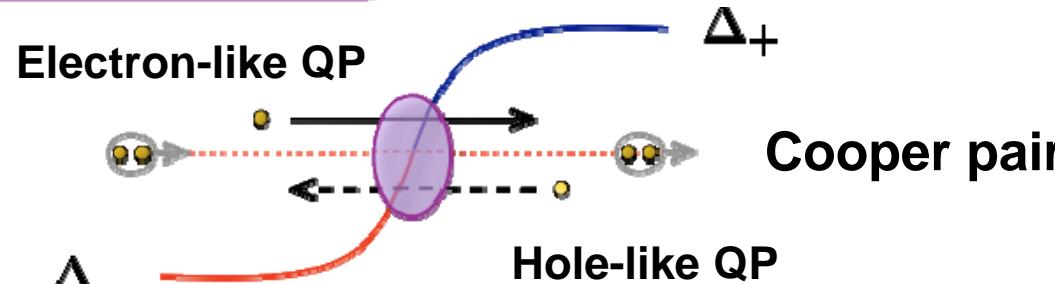


Fermi-Dirac statistics

The interface-induced state (pair amplitude) should be odd in frequency where the bulk pair potential has an **even**-frequency component since there is **no spin flip** at the interface.

Andreev bound states in inhomogeneous systems are manifestations of odd-frequency pairing amplitude

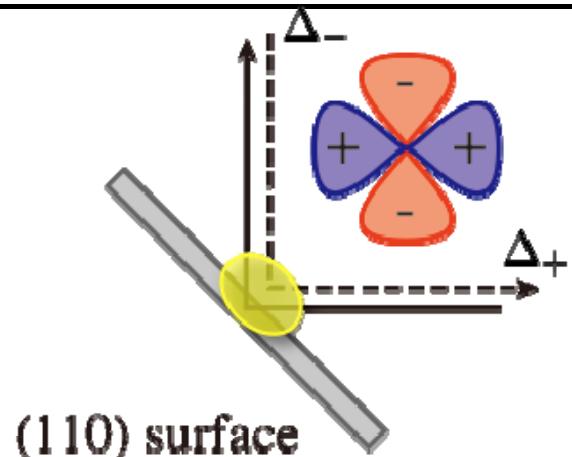
Andreev bound states



Positive pair potential

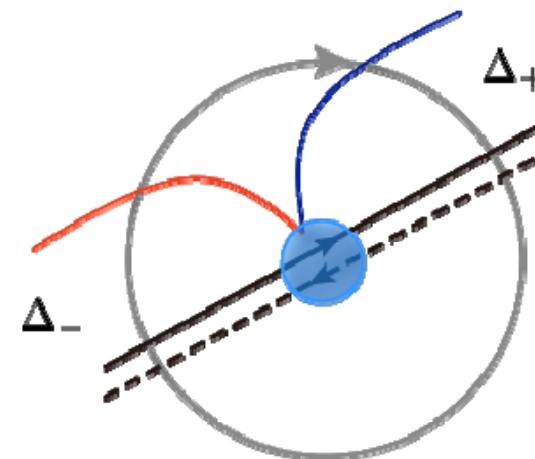
Negative pair potential

Surface: Tanaka *et al*, 2007



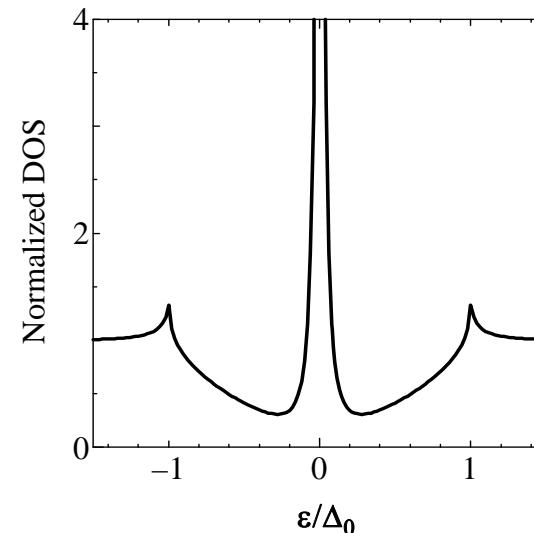
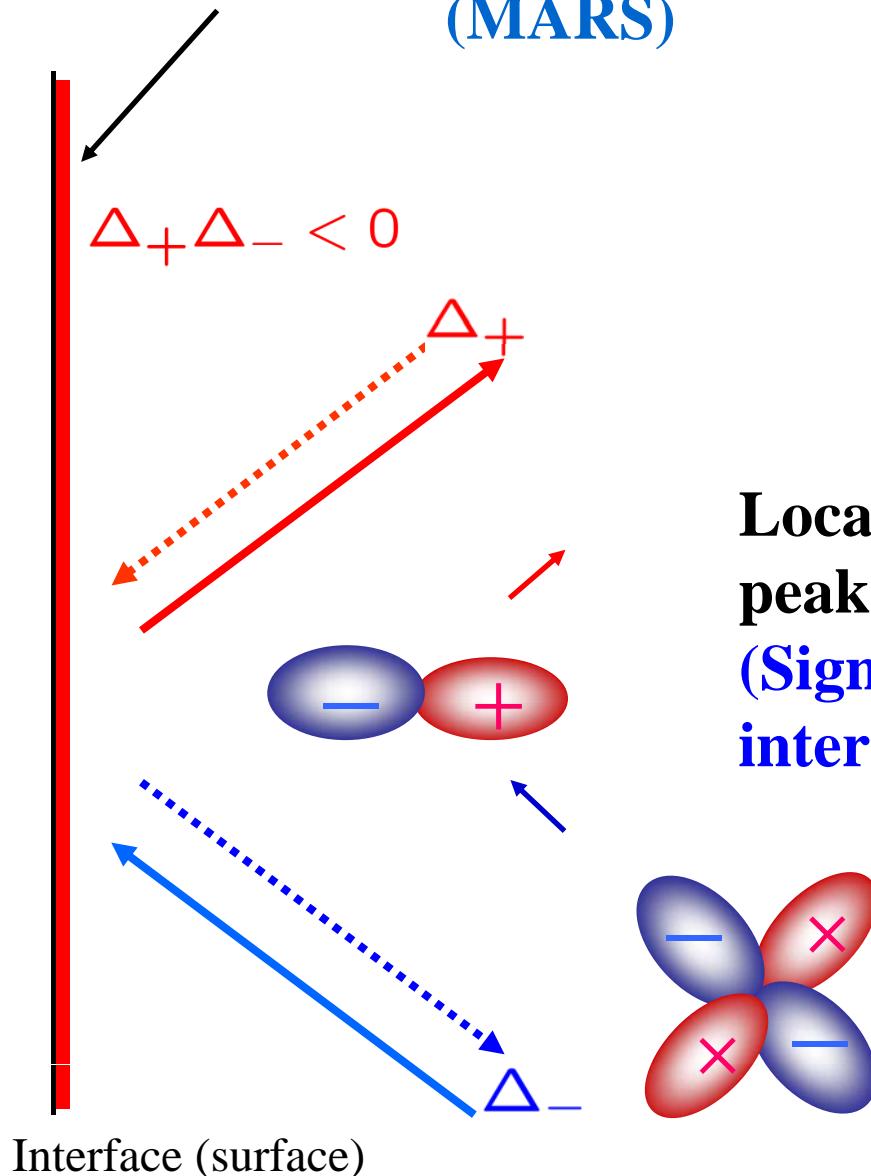
Scattering direction of QP

Vortex : Tanuma *et al*, 2009



Phase change due to a vortex

Mid gap Andreev resonant (bound) state (MARS)



Local density of state has a zero energy peak.
(Sign change of the pair potential at the interface)

Tanaka Kashiwaya PRL 74 3451 (1995),
Rep. Prog. Phys. 63 1641 (2000)
Buchholz(1981) Hara Nagai(1986)
Hu(1994) Matsumoto Shiba(1995)
Ohashi Takada(1995)
Hatsugai and Ryu (2002)

Superconducting Materials where MARS is observed

$\text{YBa}_2\text{CuO}_{7-\delta}$ (Geerk, Kashiwaya, Iguchi, Greene, Yeh,Wei..)

$\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_y$ (Ng, Suzuki, Greene....)

$\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$ (Iguchi)

$\text{La}_{2-x}\text{Ce}_x\text{CuO}_4$ (Cheska)

$\text{Pr}_{2-x}\text{Ce}_x\text{CuO}_4$ (R.L.Greene)

Sr_2RuO_4 (Mao, Meno, Kawamura,Laube)

$\kappa-(\text{BEDT-TTF})_2X$, $X=\text{Cu}[\text{N}(\text{CN})_2]\text{Br}$ (Ichimura)

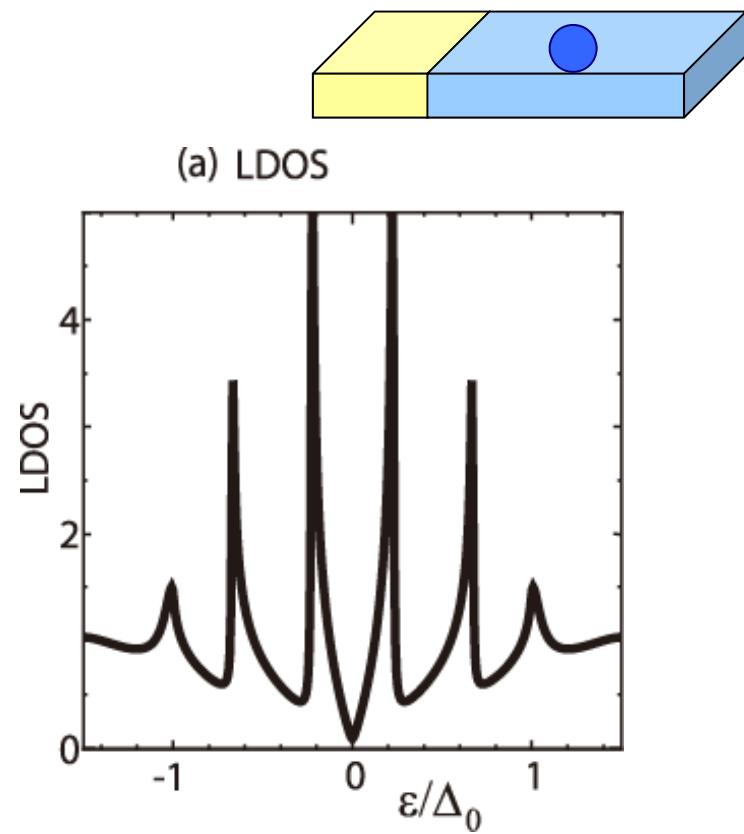
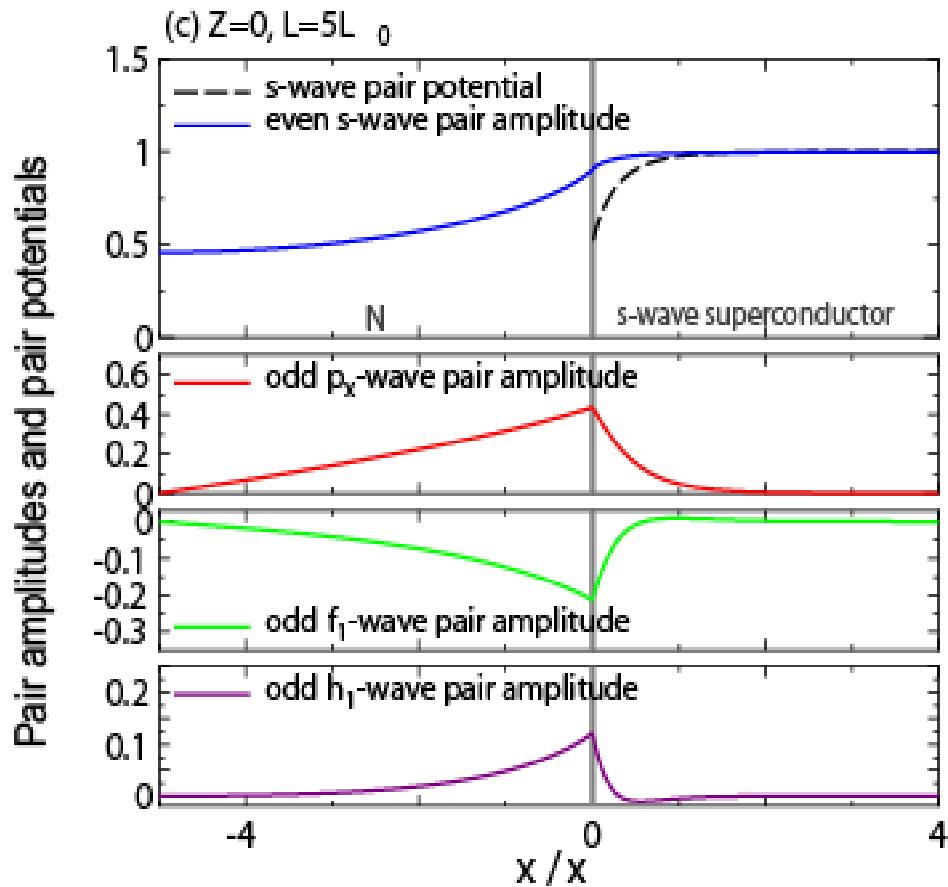
UBe_{13} (Ott)

CeCoIn_5 (Wei Greene)

$\text{PrOs}_4\text{Sb}_{12}$ (Wei)

Superfluid ^3He (Okuda, Nomura, Higashitani, Nagai)

Odd-frequency pairing state in N/S junctions (N finite length)



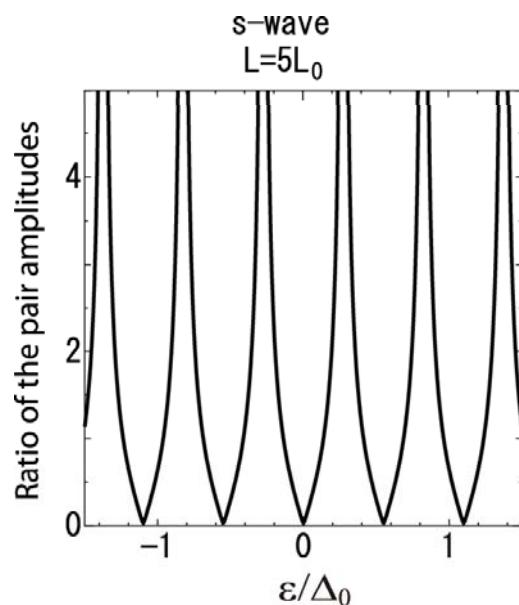
Bounds state are formed in the normal metal

Y. Tanaka, Y. Tanuma and A.A.Golubov, Phys. Rev. B **76**, 054522 (2007)

Ratio of the pair amplitude in the N region (odd/even)

$$\left| \frac{f_{1+}^{(N)}(\varepsilon, \theta)}{f_{2+}^{(N)}(\varepsilon, \theta)} \right| = \left| \tan \left(\frac{2\varepsilon}{v_F x} (L + x) \right) \right|.$$

At some energy, odd-frequency component can exceed over even frequency one.



$$\begin{aligned} f_{1+}^{(N)}(\varepsilon, \theta) & \quad \text{Odd-frequency pairing} \\ \theta = 0 & \\ x = 0 & \quad f_{2+}^{(N)}(\varepsilon, \theta) \\ & \quad \text{Even-frequency pairing} \end{aligned}$$

Hidden odd-frequency component in the s-wave superconductor junctions

Ratio of the pair amplitude at the N/S interface and the bound state level

$\Delta_0 \gg \varepsilon$ **Bound states condition (Z=0)**
(McMillan Thomas Rowell)

$$\varepsilon_n = \frac{\pi v_F x}{2L} (n + 1/2), \quad n = 0, 1, 2, \dots$$

$$\left| \frac{f_{1+}^{(N)}(\varepsilon, \theta)}{f_{2+}^{(N)}(\varepsilon, \theta)} \right| = |\tan(\pi/2 + \pi n)| = \infty.$$

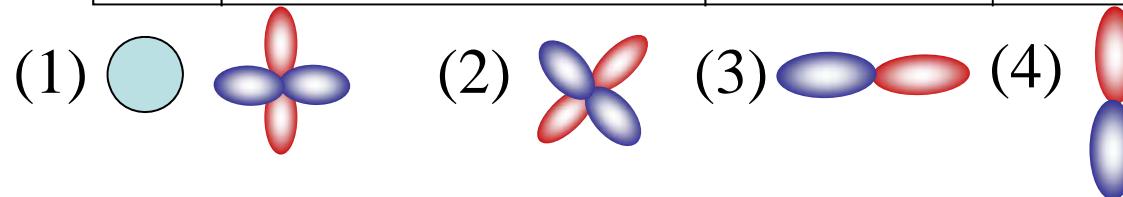
$f_{1+}^{(N)}(\varepsilon, \theta)$
Odd-frequency pairing

$f_{2+}^{(N)}(\varepsilon, \theta)$
Even-frequency pairing

Bound states are due to the generation of the odd-frequency Cooper pair amplitude

Symmetry of the Cooper pair (No spin flip)

	Bulk state	Sign change (MARS)	Interface-induced symmetry (subdominant component)
(1)	ESE ($s, d_{x^2-y^2}$ -wave)	No	ESE + (OSO)
(2)	ESE (d_{xy} -wave)	Yes	OSO + (ESE)
(3)	ETO (p_x -wave)	Yes	OTE + (ETO)
(4)	ETO (p_y -wave)	No	ETO + (OTE)



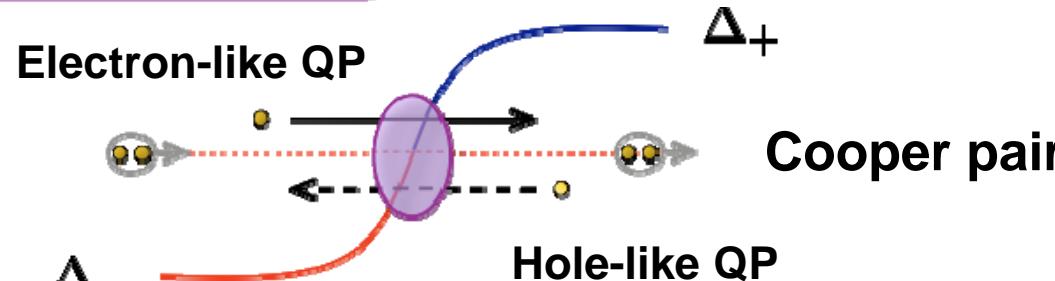
- **ESE (Even-frequency spin-singlet even-parity)**
- **ETO (Even-frequency spin-triplet odd-parity)**
- **OTE (Odd-frequency spin-triplet even-parity)**
- **OSO (Odd-frequency spin-singlet odd-parity)**

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Andreev bound states in inhomogeneous systems are manifestations of odd-frequency pairing amplitude

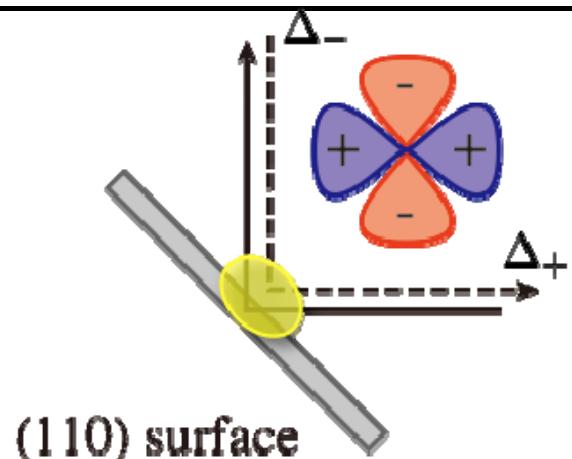
Andreev bound states



Positive pair potential

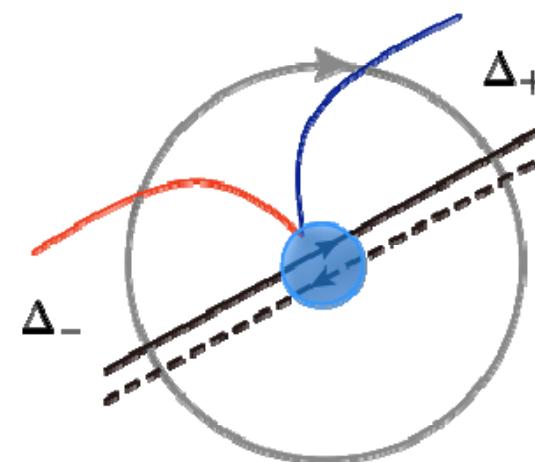
Negative pair potential

Surface: Tanaka *et al*, 2007



Scattering direction of QP

Vortex : Tanuma *et al*, 2009



Phase change due to a vortex

Symmetry of the Cooper pair in a vortex core

l ; angular momentum	m ; vorticity	bulk	Center of the vortex core
Even	Even	ESE (s-wave..)	ESE
Even	Odd	ESE (s-wave..)	OSO
Odd	Even	ETO (chiral p-wave)	ETO
Odd	Odd	ETO (chiral p-wave)	OTE

$$\Delta(\mathbf{r}) = \Delta_0 \exp(il\varphi) \tanh\left(\frac{\sqrt{x^2 + y^2}}{\xi}\right) \left(\frac{x + iy}{\sqrt{x^2 + y^2}}\right)^m$$

ESE (Even-frequency spin-singlet even-parity)

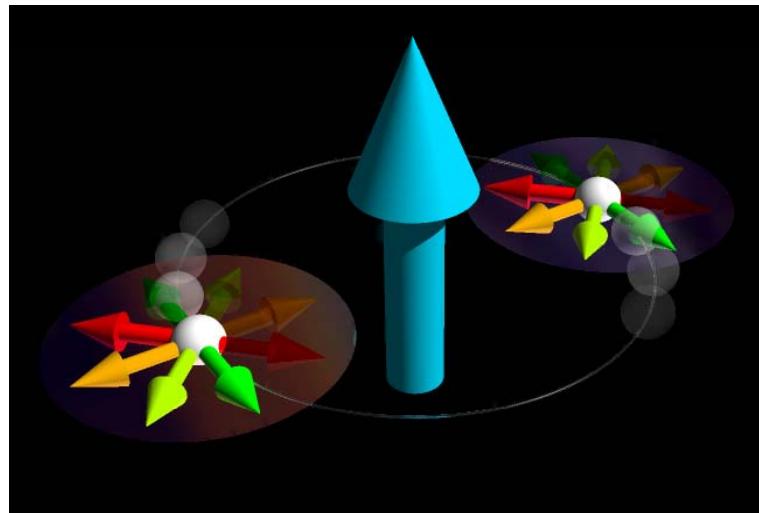
ETO (Even-frequency spin-triplet odd-parity)

OTE (Odd-frequency spin-triplet even-parity)

OSO (Odd-frequency spin-singlet odd-parity)

Yokoyama et al., Physical Review B, Vol. 78, 012508, 2008

Vortex core spectroscopy of chiral p-wave superconductors



Sr₂RuO₄
Maeno (1994)

Y. Tanuma, N. Hayashi, Y. Tanaka, A. A. Golubov

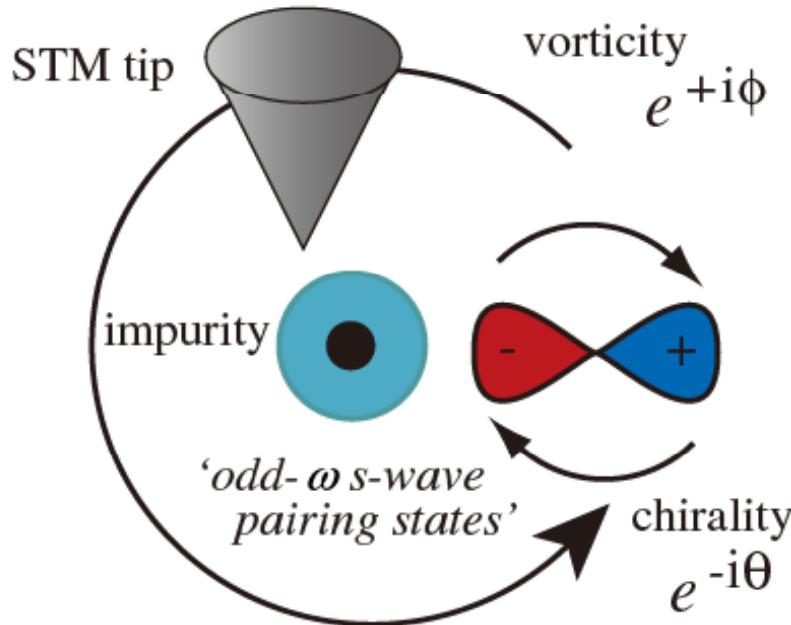
Phys. Rev. Lett. 102, 117003 (2009)

Chirality and vorticity: Y. Kato and N. Hayashi (2000, 2001, 2002)

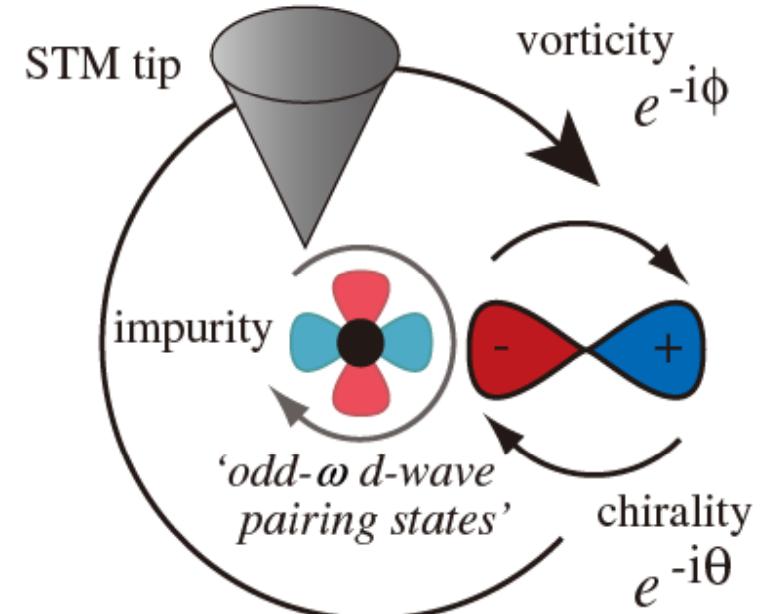
**Focus on the impurity scattering effect
(Born limit)**

Difference of the angular momentum of the odd-frequency pair at the core center

(a) Antiparallel chiral p-wave vortex



(b) Parallel chiral p-wave vortex



Angular momentum at the center of core; $\textcolor{red}{l+m}$

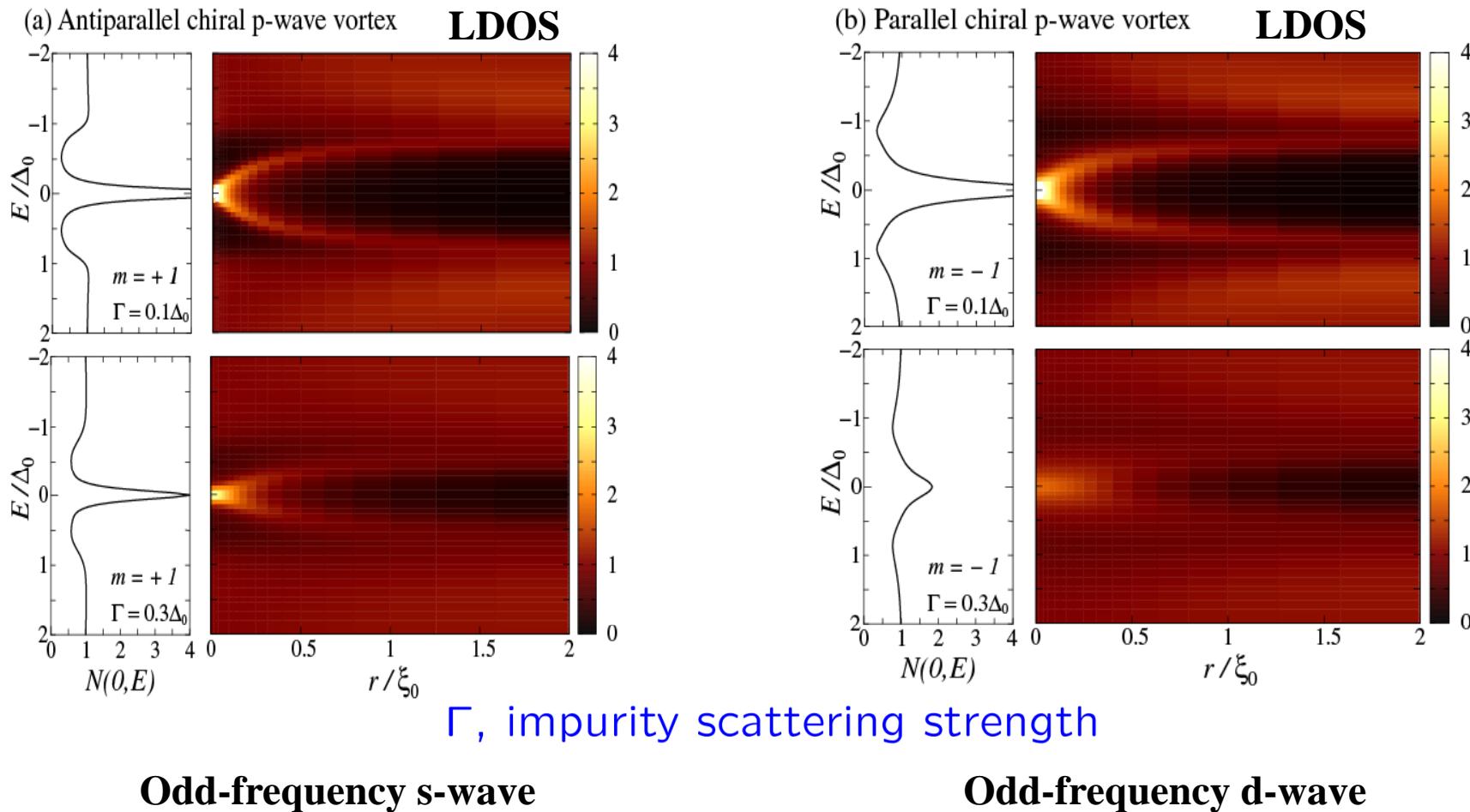
l : angular momentum

m : vorticity

Tanuma, Hayashi, Tanaka Golubov Phys. Rev. Lett. 102, 117003 (2009).

Chirality and vorticity: Y. Kato and N. Hayashi (2000, 2001, 2002)

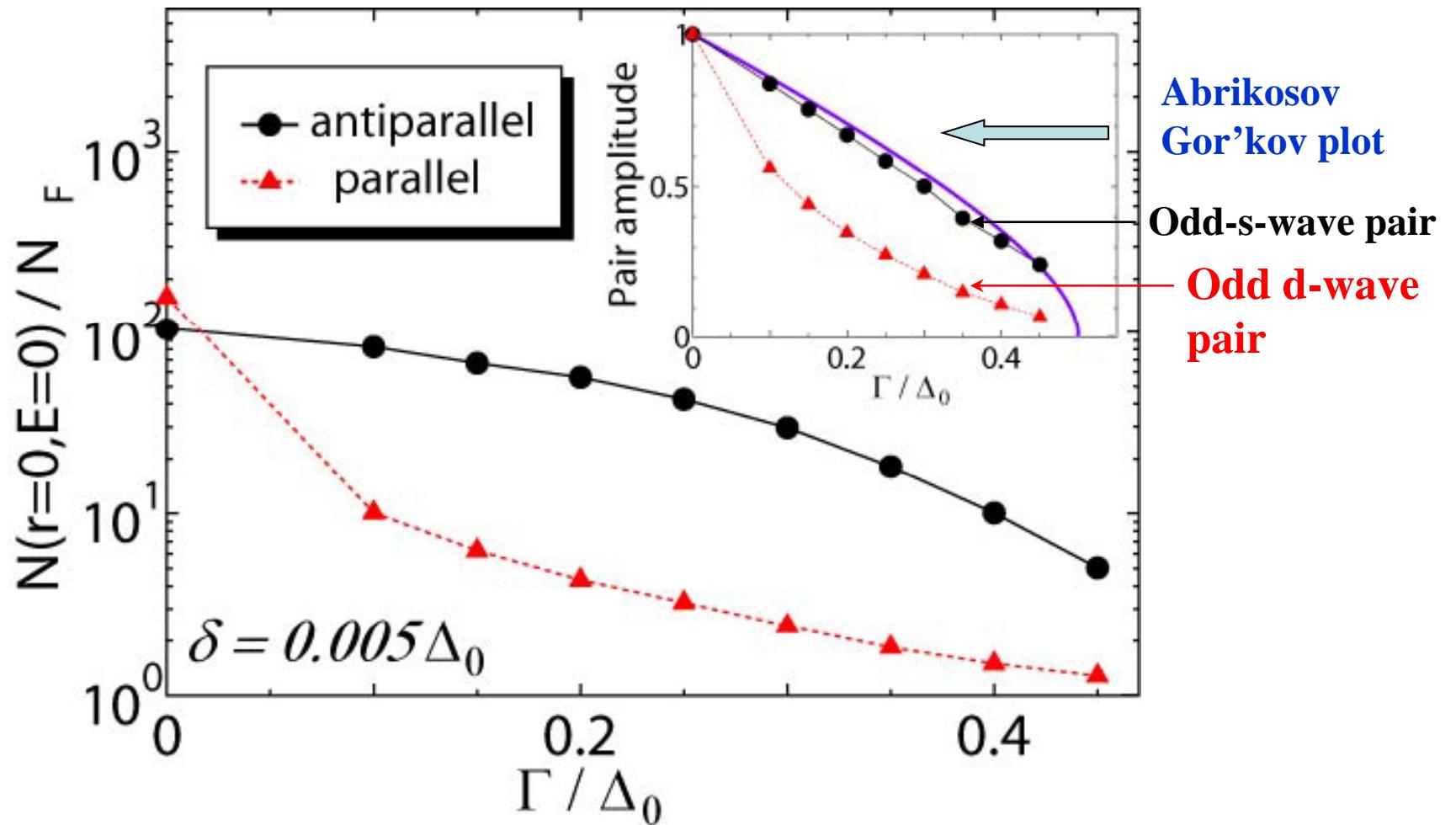
Impurity effect (Born approximation)



Odd frequency s-wave state; more robust against the impurity scattering Tanuma, Hayashi, Tanaka Golubov, PRL 102 117003 (2009).

Chirality and vorticity: Y. Kato and N. Hayashi (2000, 2001, 2002)

Impurity scattering



Chirality and vorticity: Y. Kato and N. Hayashi (2000, 2001, 2002)

Chiral domain and vortex

Parallel vortex	Antiparallel vortex	Parallel vortex	Antiparallel vortex
			
Odd d-wave	Odd s-wave	Odd d-wave	Odd s-wave
$p_x + ip_y$	$p_x - ip_y$	$p_x + ip_y$	$p_x - ip_y$
Weak ZEP	Strong ZEP	Weak ZEP	Strong ZEP

We can detect the presence of chiral domain by vortex spectroscopy via odd-frequency Cooper pair.

Phys. Rev. Lett. 102 117003 (2009).

Summary (vortices)

(1) If we consider Abrikosov vortex ($m=1$), only the odd-frequency Cooper pair is possible at the center of the vortex core.

Physical Review B, Vol. 78, 012508, 2008

(2) Vortex core spectroscopy in chiral p-wave superconductor in the presence of impurity enables us to identify the presence of chirality and the odd-frequency pairing.

Phys. Rev. Lett. 102, 117003 (2009).

Summary

- (1) Ubiquitous presence of the odd-frequency pairs in inhomogeneous systems.
- (2) Low energy Andreev bound states can be expressed in terms of odd-frequency pairing (proximity effect and vortices).
- (3) Odd-frequency Cooper pairing is realized at the center of a vortex core -> allows to identify the presence of chirality