

Nuclear Spin Ordering in Solid ^3He

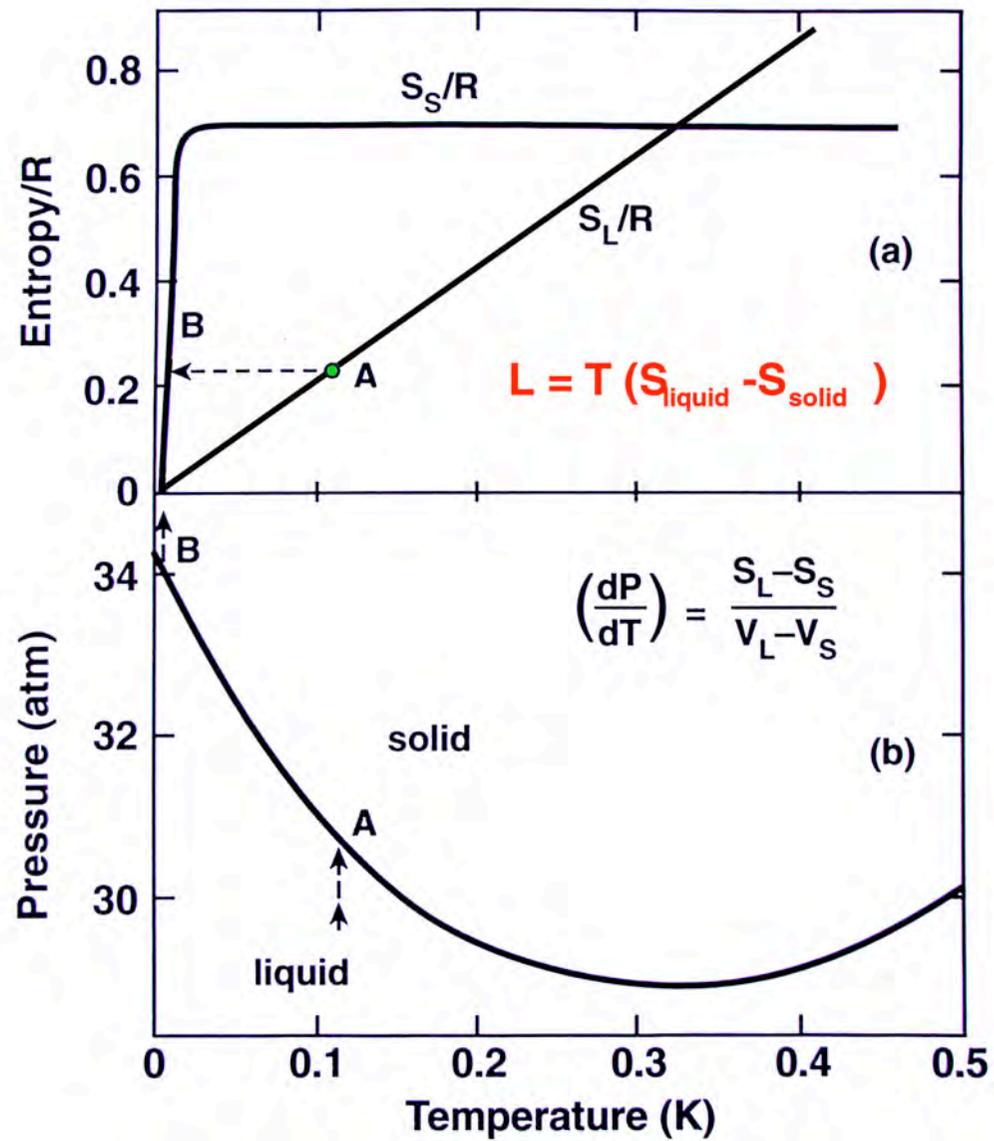


Douglas Osheroff
Stanford University

PSM 2010 Symposium

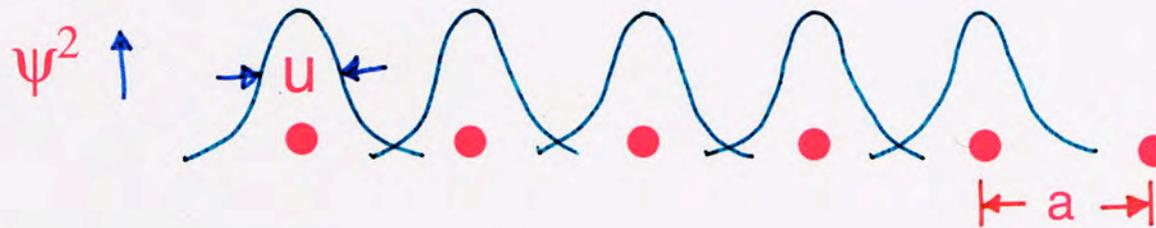
Yokohama, Japan
12 March 2010

Pomeranchuk's Conjecture: 1950



Fermi-Particle Exchange Spin Interactions:

a) Large zero point energy:



$$u = a/3$$

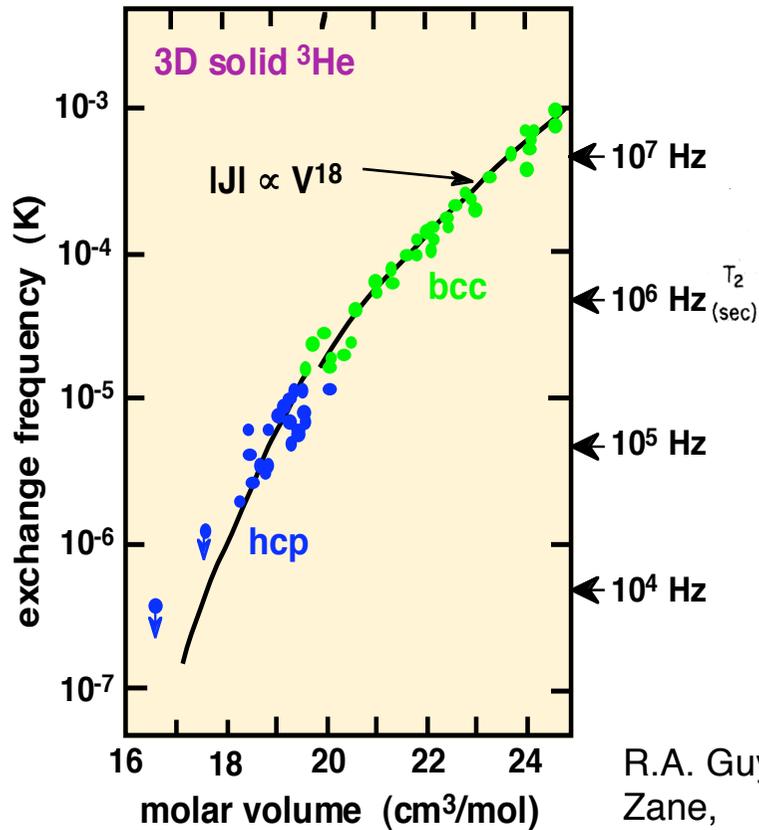
b) Atoms exchange lattice sites once in every $\sim 10^4$ zero point oscillations.

c) Lowest energy for two-particle exchange

- Symmetric orbital wave function
- Antisymmetric spin wave function

Early studies of exchange interactions in solid ^3He (1960s)

Volume dependence of J



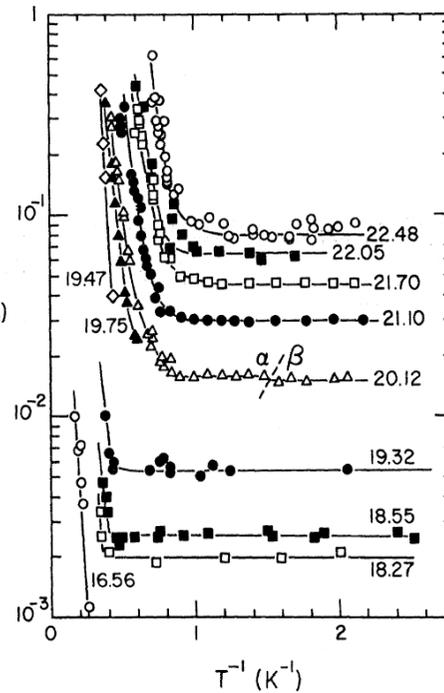
R.A. Guyer, R.C. Richardson and L.I. Zane,
Rev. Mod. Phys. **43**, 532 (1971)

strong volume dependence

$$|J| \propto V^{18}$$

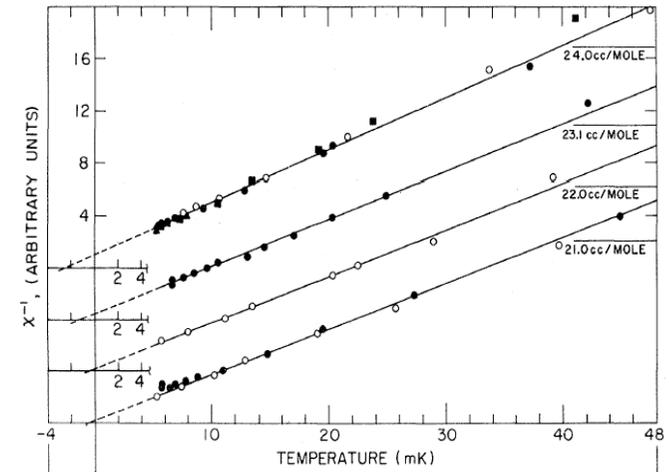
$J \approx -1$ mK at melting
pressure

T_2 in NMR



exchange
narrowing
 $T_2 \propto |J|$

Magnetic susceptibility



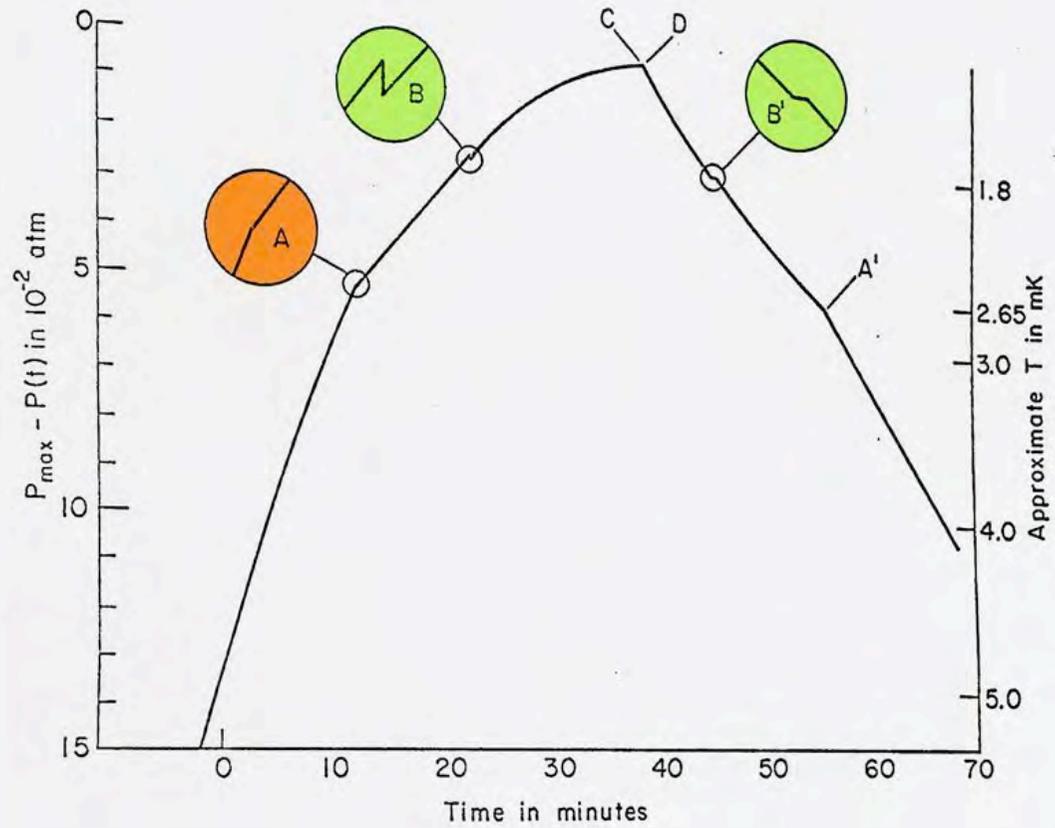
W.P. Kirk et al., PRL **23**, 833
(1969)

antiferromagnetic

$$\theta_W (= -4J) < 0$$

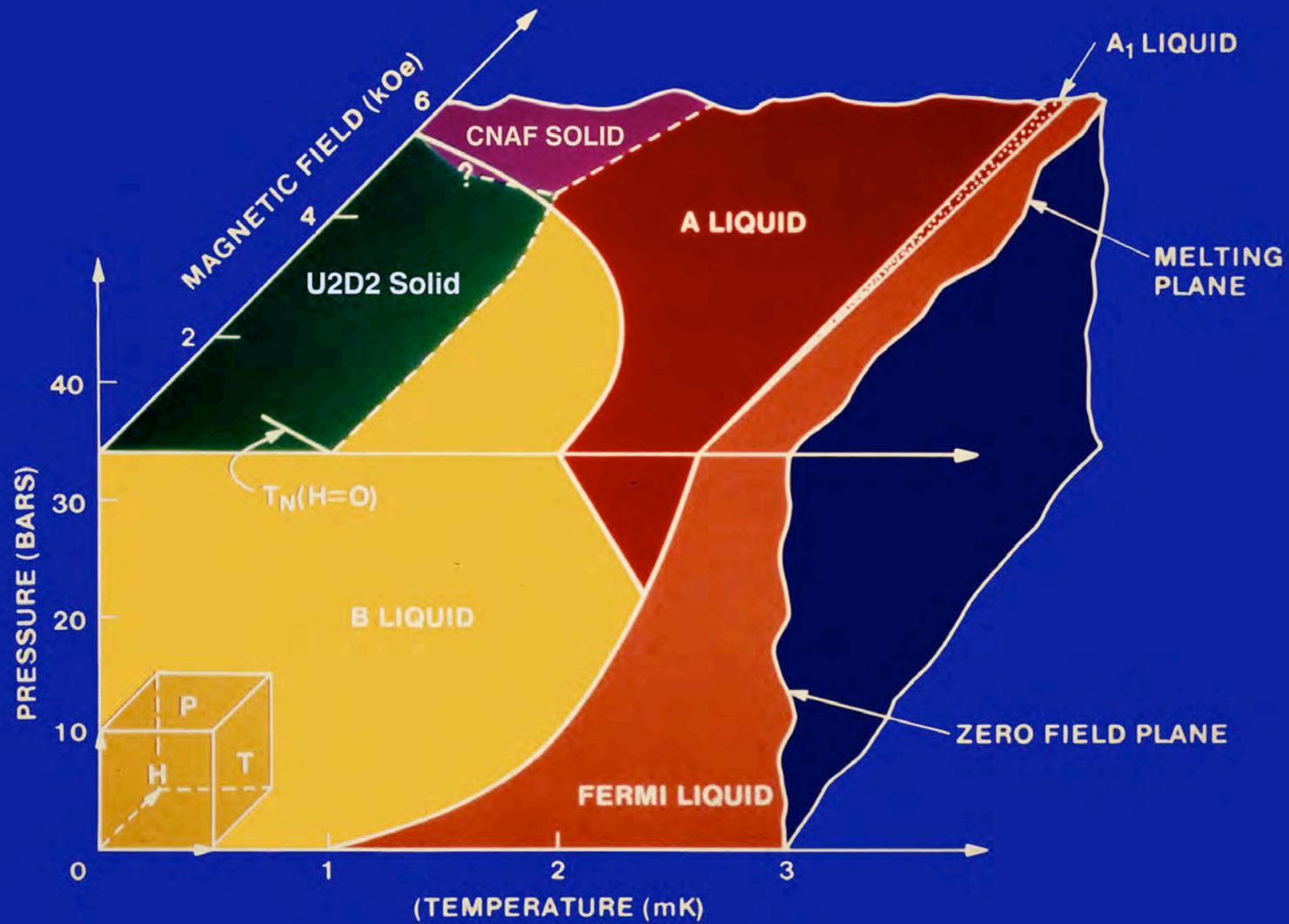
Thanks to Hiroshi Fukuyama

Full Pressurization Curve

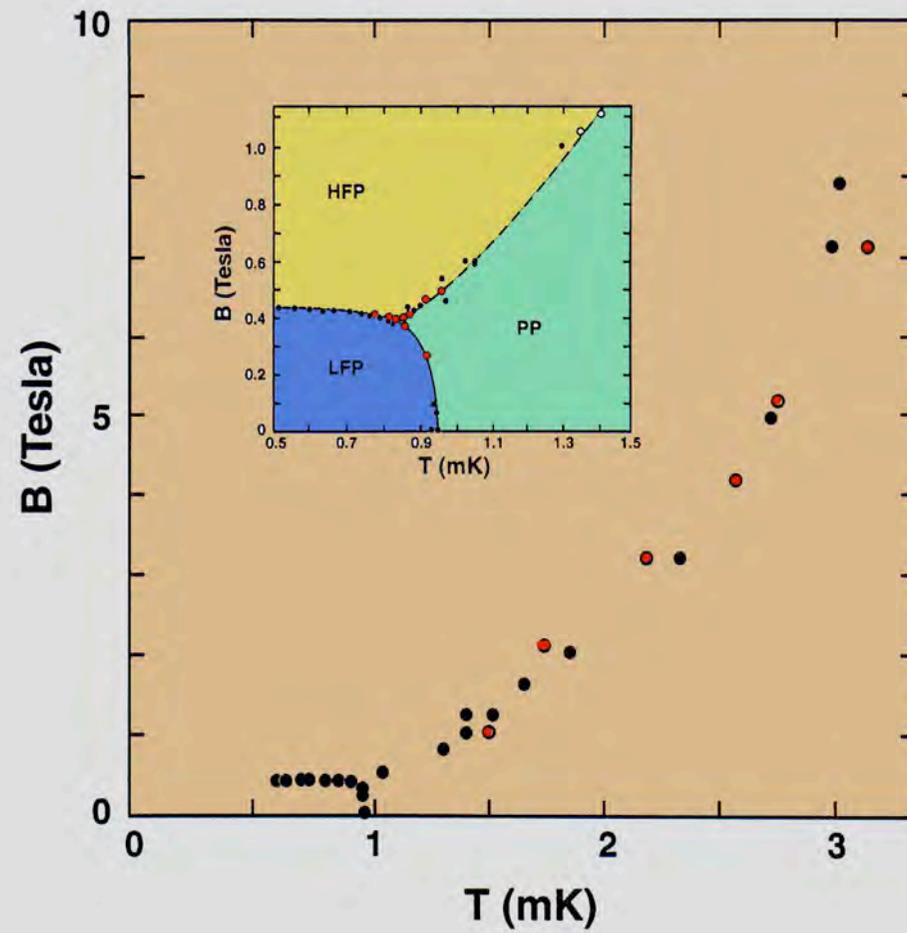


$$\Delta P \sim (1/\Delta V) \int_{T_i}^{T_f} S_{\text{solid}}(T) dT$$

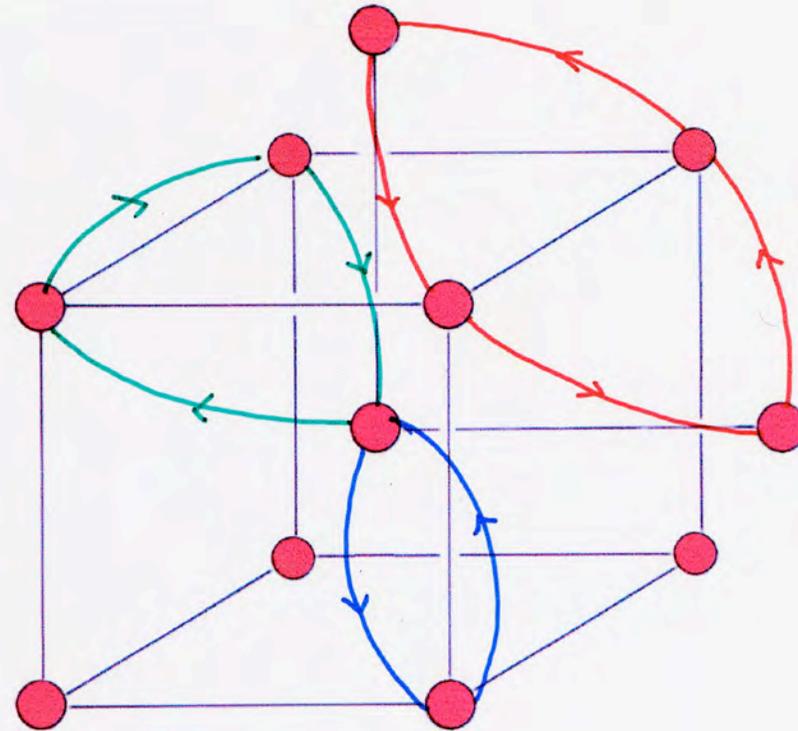
Ordered Helium Three Phase Diagram



Helium Three Magnetic Phase Diagram



Multiple Exchange Hamiltonian

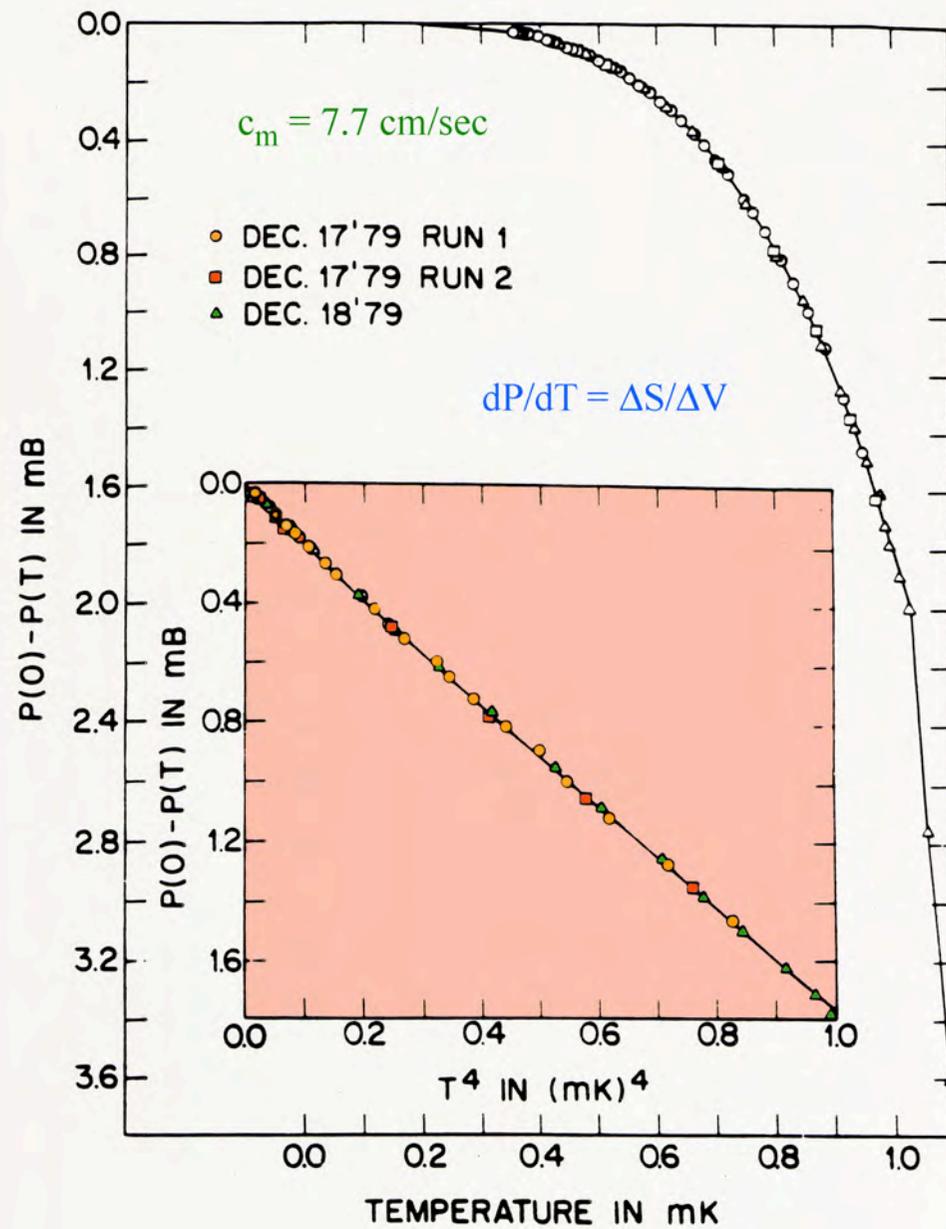


Even # exchanges antiferromagnetic,
odd # exchanges promote ferromagnetism.

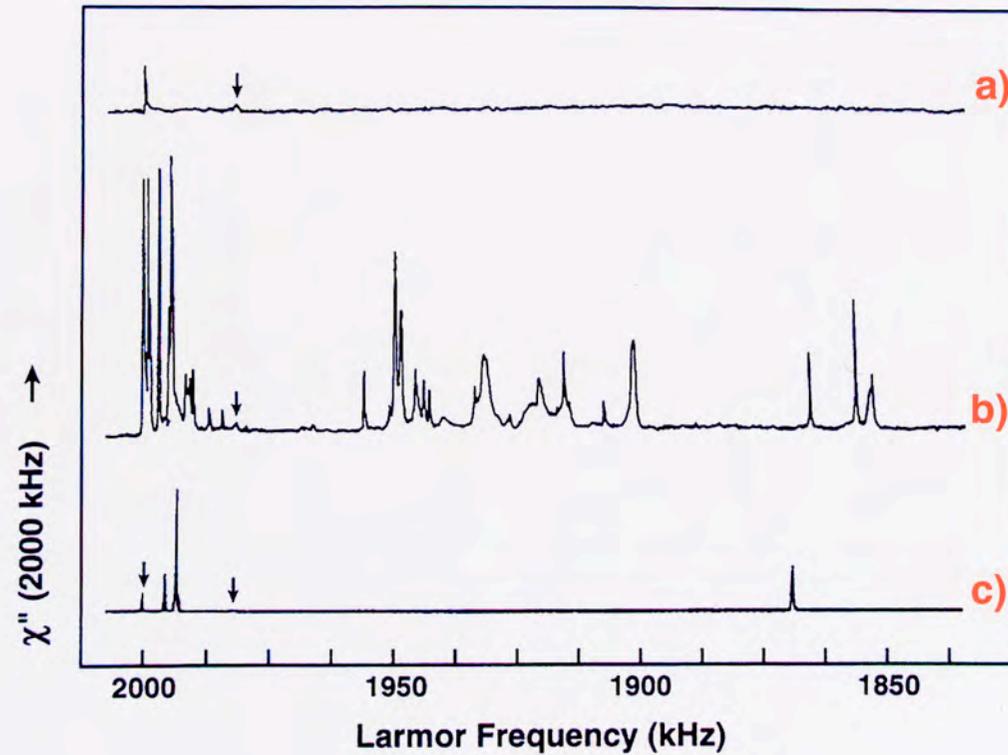
— J_{nn} — J_t — K_p

Frustration

Osheroff and Yu, Phys. Letts. 77A, 458 (1980)



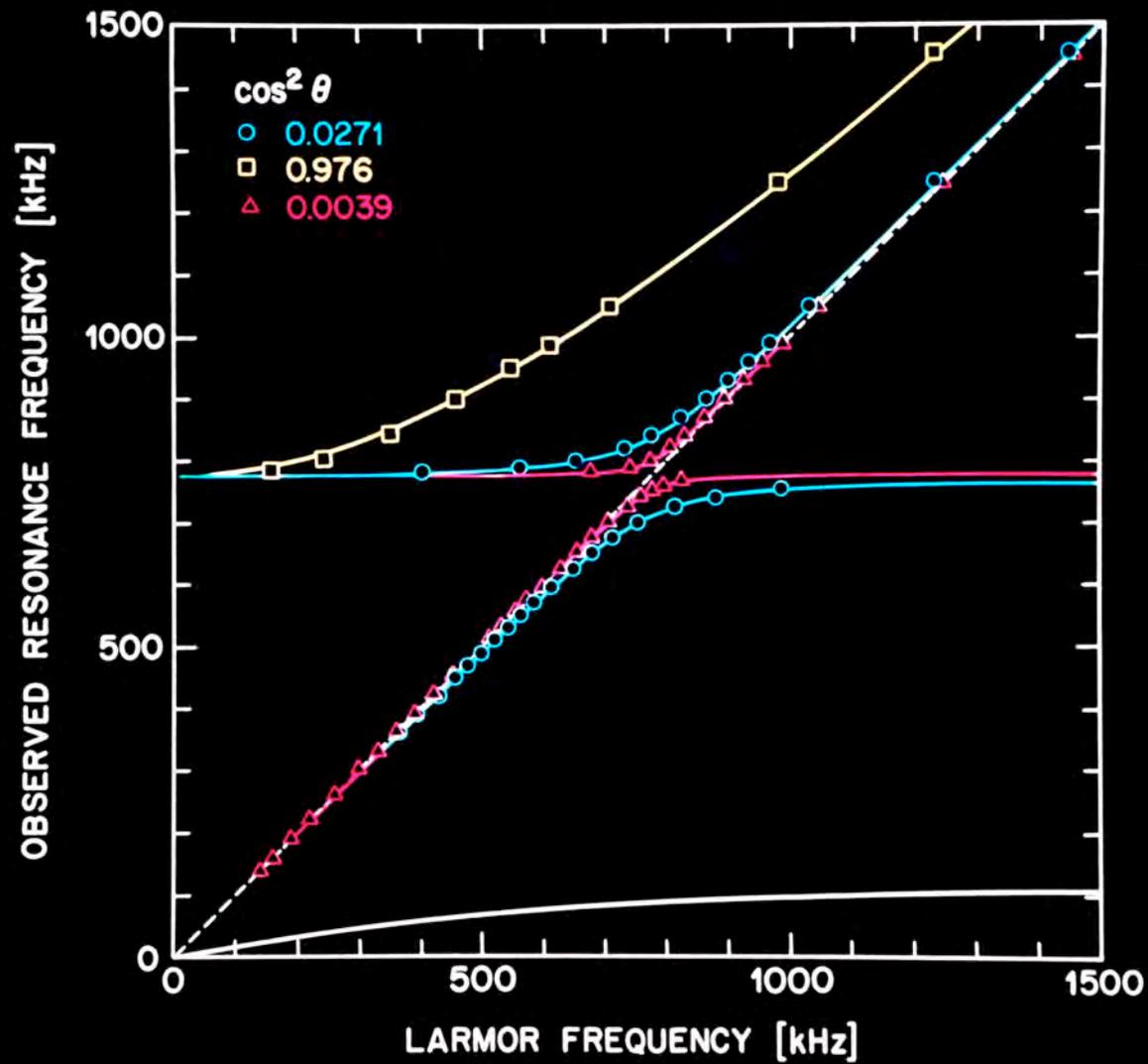
Growth of Single Crystals of Solid ^3He



a) All-liquid NMR signal

b) Heat pulse growth pattern

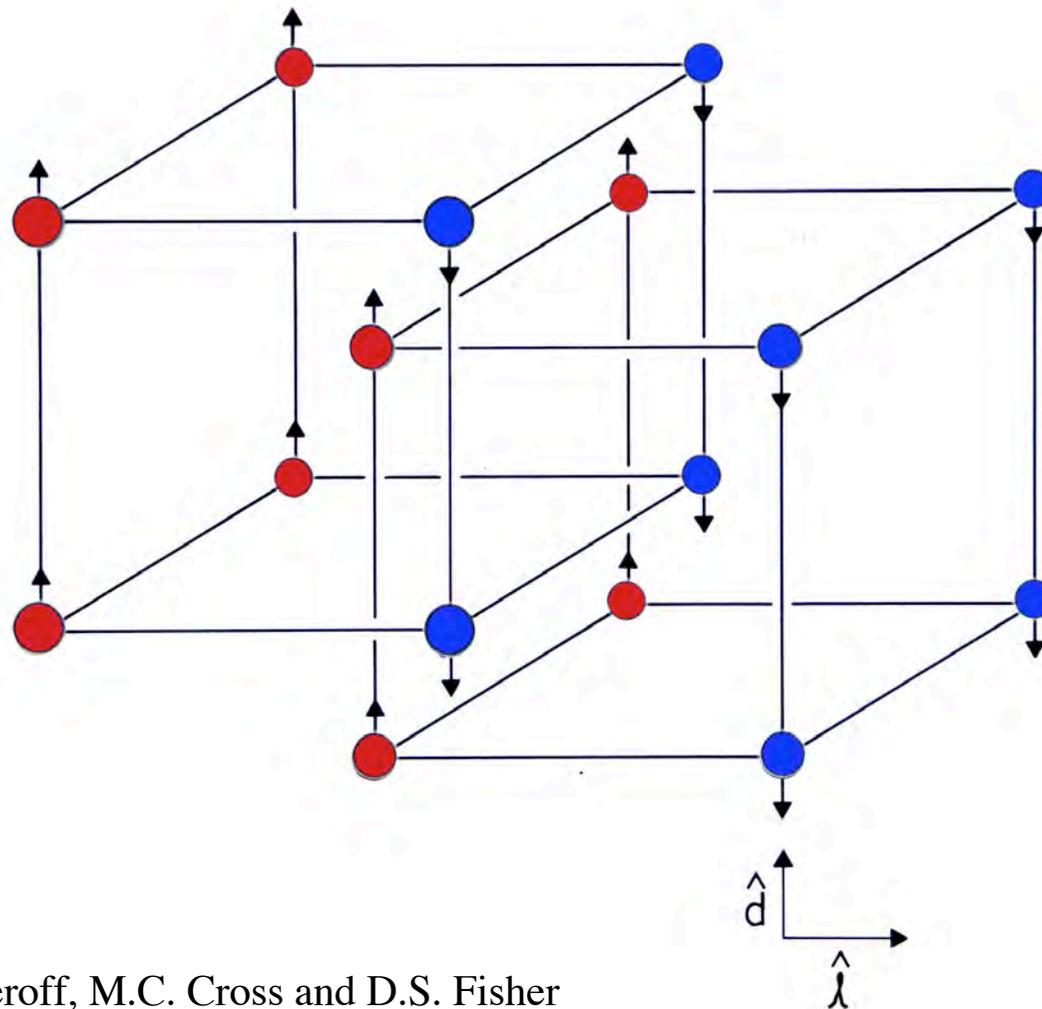
c) Seed XI re-growth pattern



$$\left(\nu_i^{\pm}\right)^2 = \frac{1}{2} \left\{ \nu_L^2 + \Omega_0^2 \pm \sqrt{(\nu_L^2 - \Omega_0^2)^2 + 4\nu_L^2 \Omega_0^2 \cos^2 \Gamma_i} \right\}$$

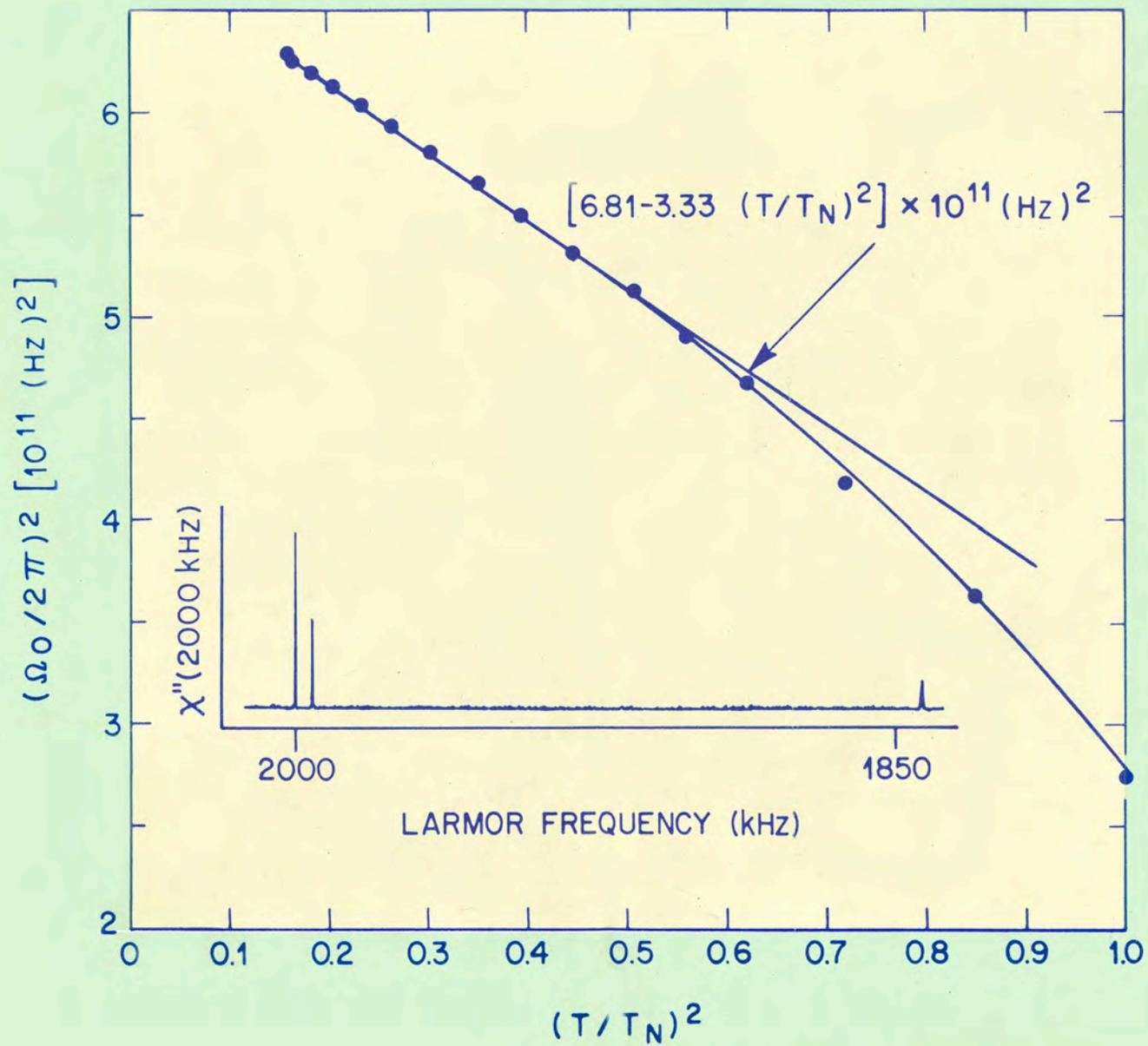
The U2D2 Phase of Solid ^3He

(after Star Wars' R2D2)



D.D. Osheroff, M.C. Cross and D.S. Fisher
Phys. Rev. Letts. 44, 792 (1980).





$$\Omega_0 = \left(\frac{\sum_{i=1}^3 v_i^4 - v_L^2 \sum_{i=1}^3 v_i^2}{\sum_{i=1}^3 v_i^2 - 2 v_L^2} \right)^{1/2}$$

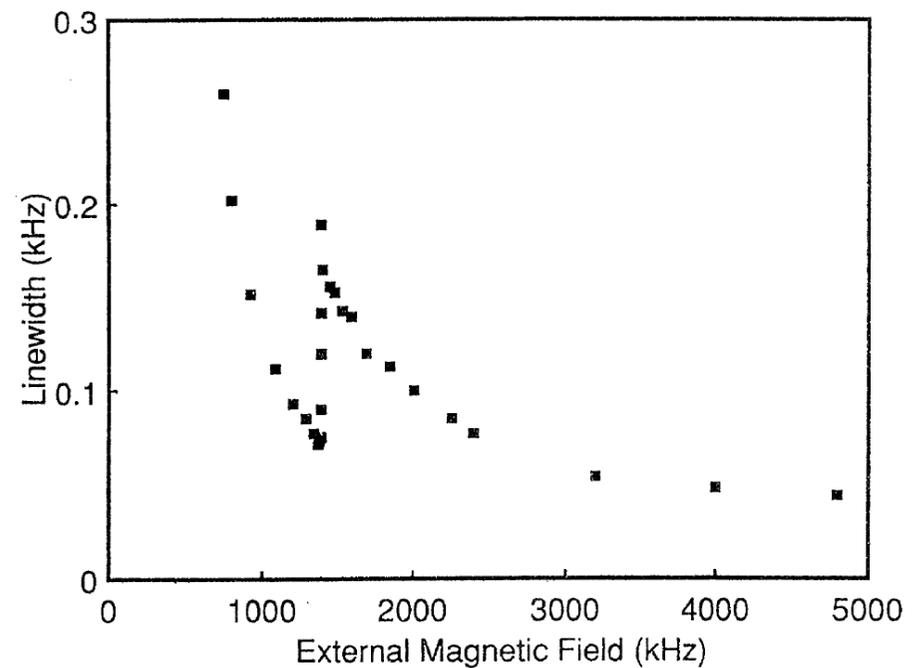
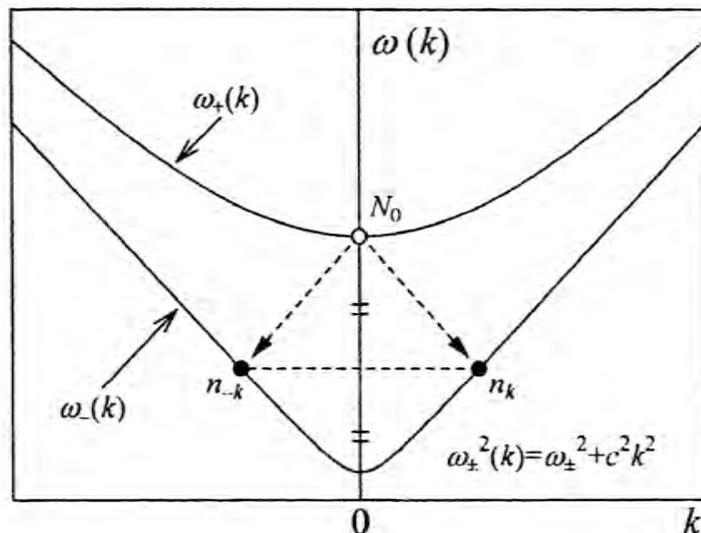
$$\cos^2 \Gamma_i = \frac{(v_i^2 - \Omega_0^2)(v_i^2 - v_L^2)}{\Omega_0^2 v_L^2}$$

$$\Omega_0 = \left(\frac{v_i^2 (v_i^2 - v_L^2)}{v_i^2 - v_L^2 (1 - \cos^2 \Gamma_i)} \right)^{1/2}$$

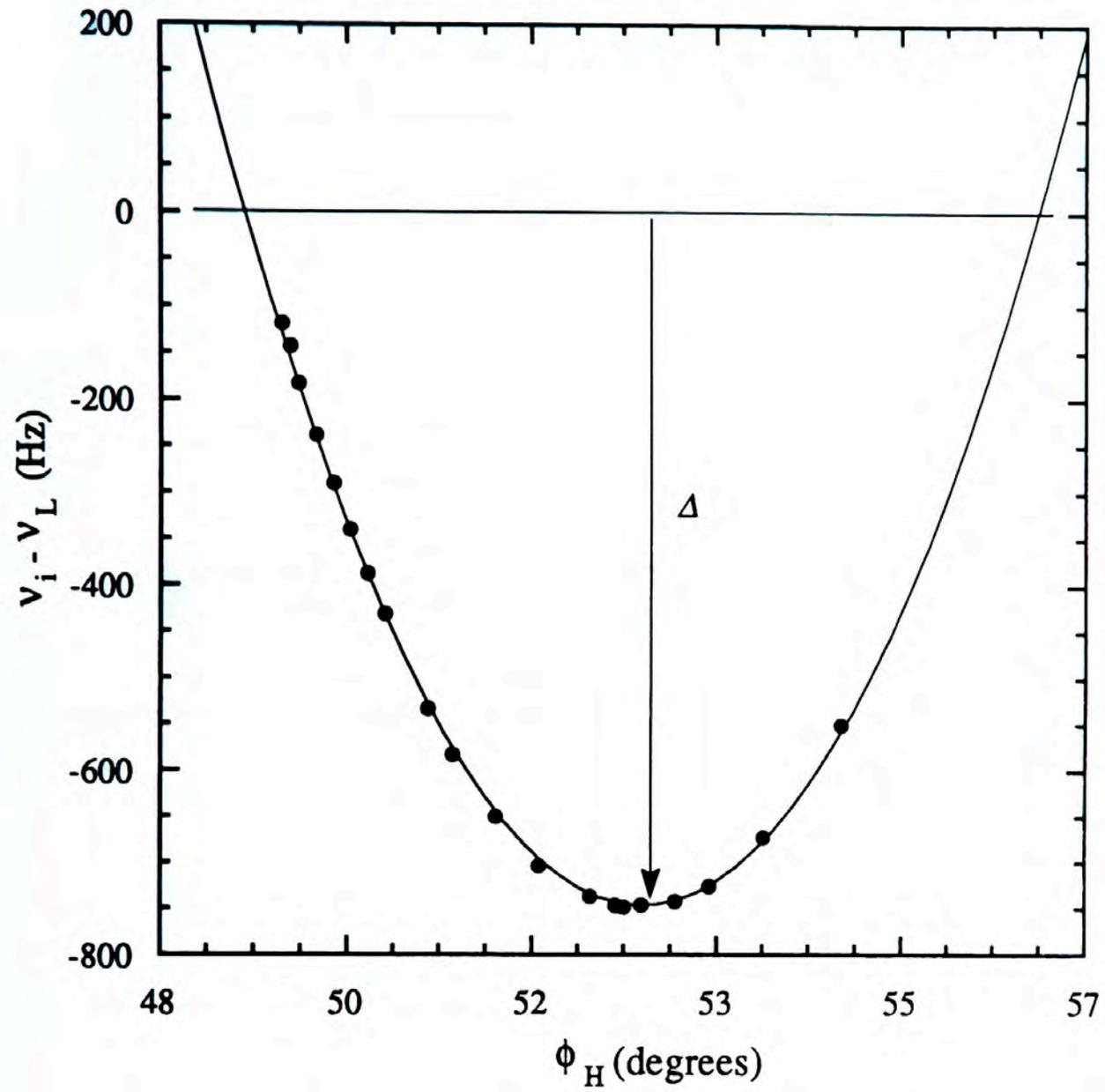
Three Magnon Decay Processes

- A magnon from the upper branch ($\omega > \gamma H$) can decay into two magnons in the lower branch provided energy and quasi-momentum can be conserved.
- For uniform spin precession in the upper branch, this loss is manifested by an abrupt broadening of the NMR line as shown.

$$\begin{cases} \dot{\mathbf{S}} = \gamma \mathbf{S} \times \mathbf{H} - \lambda (\mathbf{d} \cdot \mathbf{l})(\mathbf{d} \times \mathbf{l}) \\ \dot{\mathbf{d}} = \mathbf{d} \times \left(\gamma \mathbf{H} - \frac{\gamma^2}{\chi_{\perp}} \mathbf{S} \right) \end{cases}$$



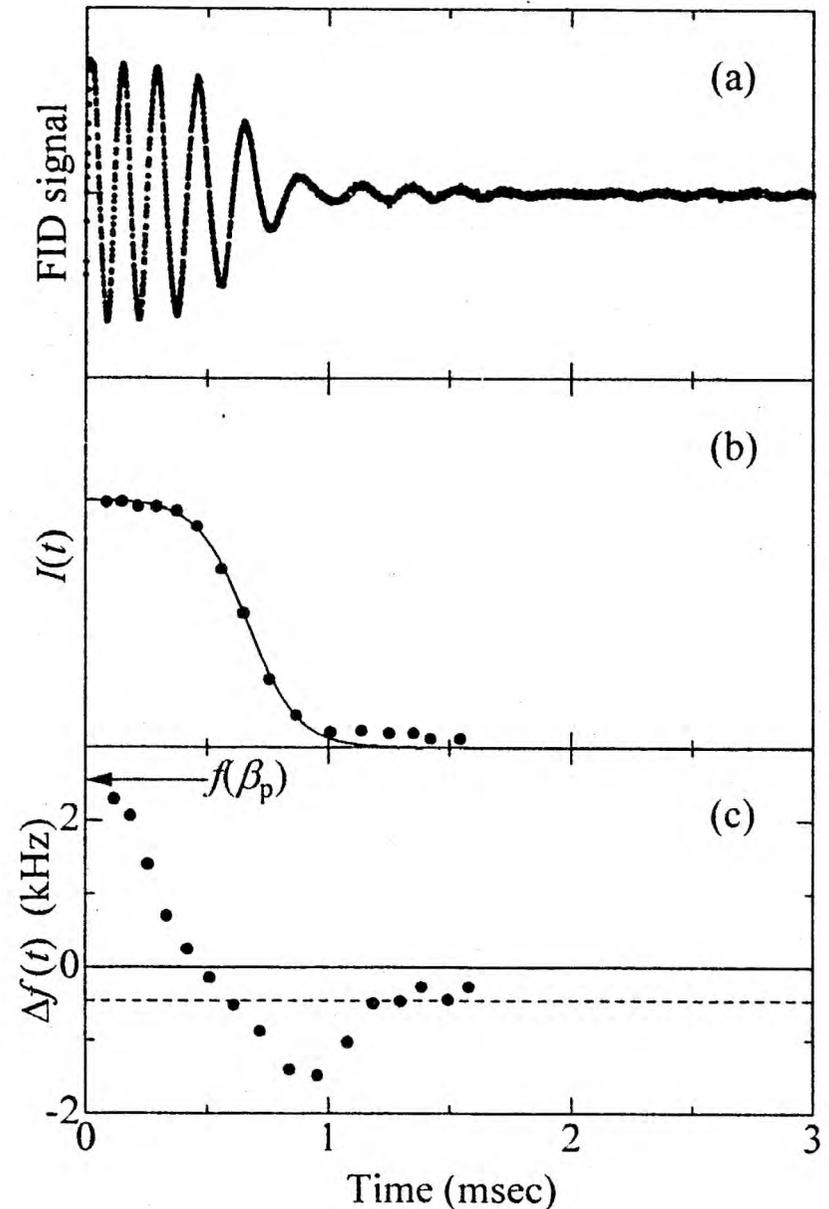
Thanks to Takao Mizusaki



Three Magnon Decay Processes

- Rapid three magnon decay from the upper frequency mode due to stimulated emission.
- Ohmi and Fomin showed that population of lower magnon mode produces a negative frequency shift in the upper mode.
- Non-linear spin dynamics in the U2D2 phase are complex but based on simple Hamiltonian.

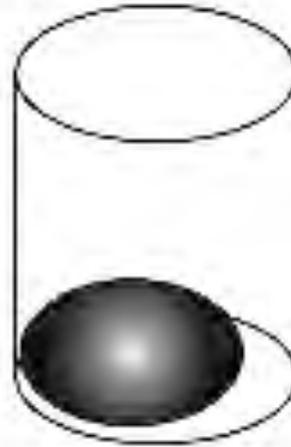
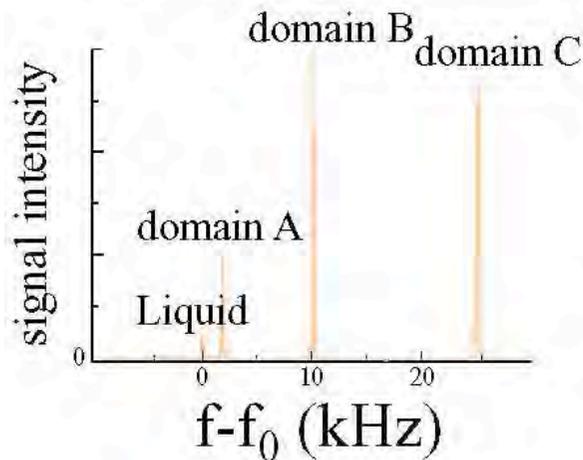
Mizusaki et al: J. Low Temp. Phys. **89** 365-373 (1992)



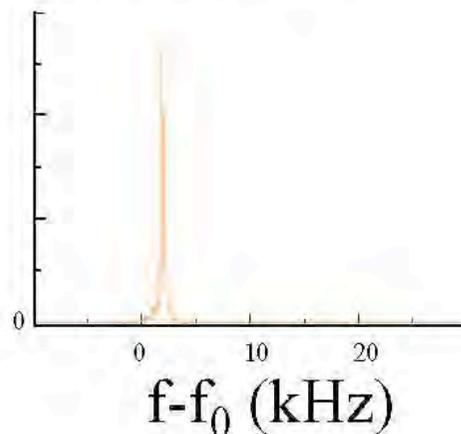
Domain Memory Between U2D2 and CNAF Phases



U2D2 spectrum



HFP spectrum



- Grow a 3-domain U2D2 sample near lower critical field.

- Measure the intensity from the three domains.

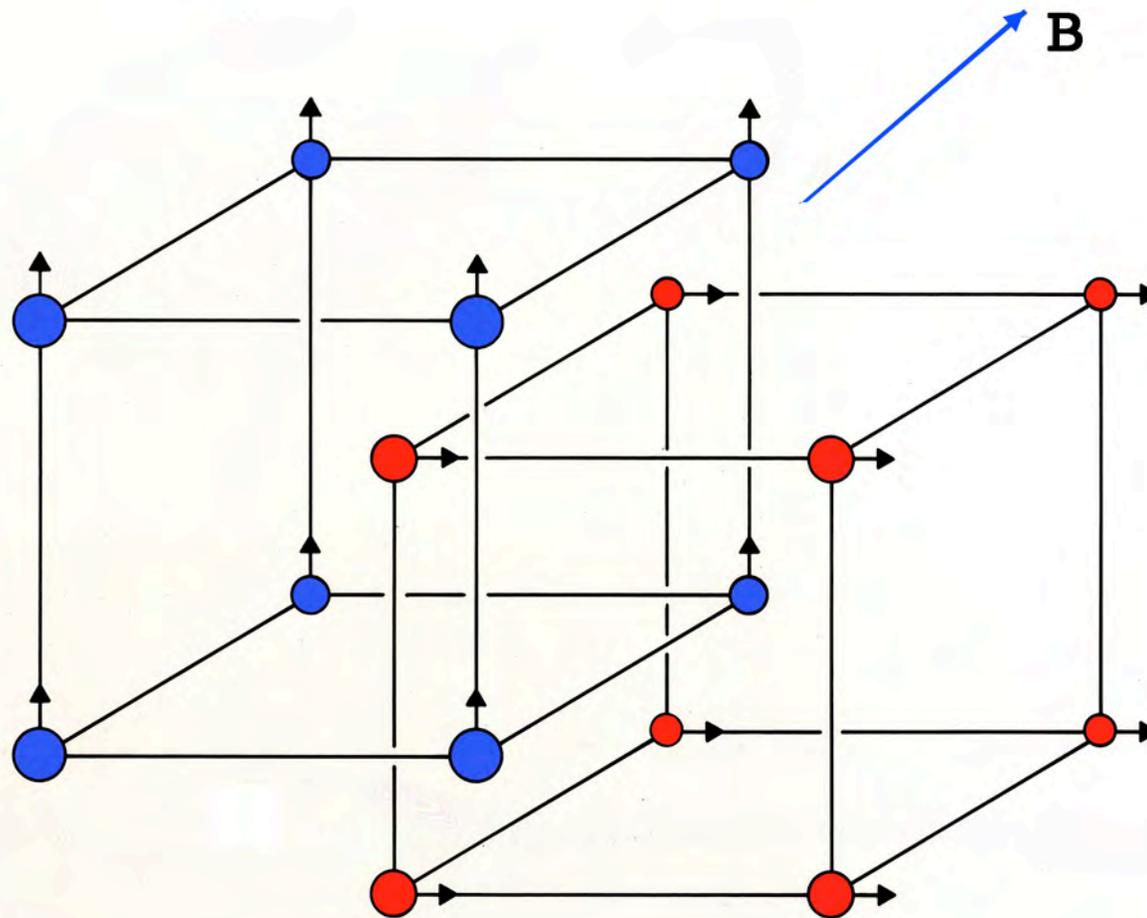
- Slowly raise magnetic field into the CNAF field Region, wait, and then lower the field back.

- Measure the intensity in the three domains again.

- For constrained xls domains are the same. For unconstrained samples, they are very different.

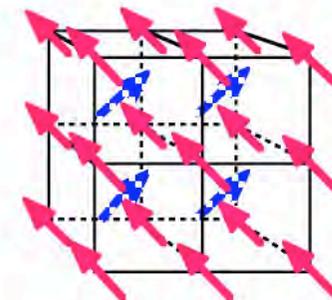
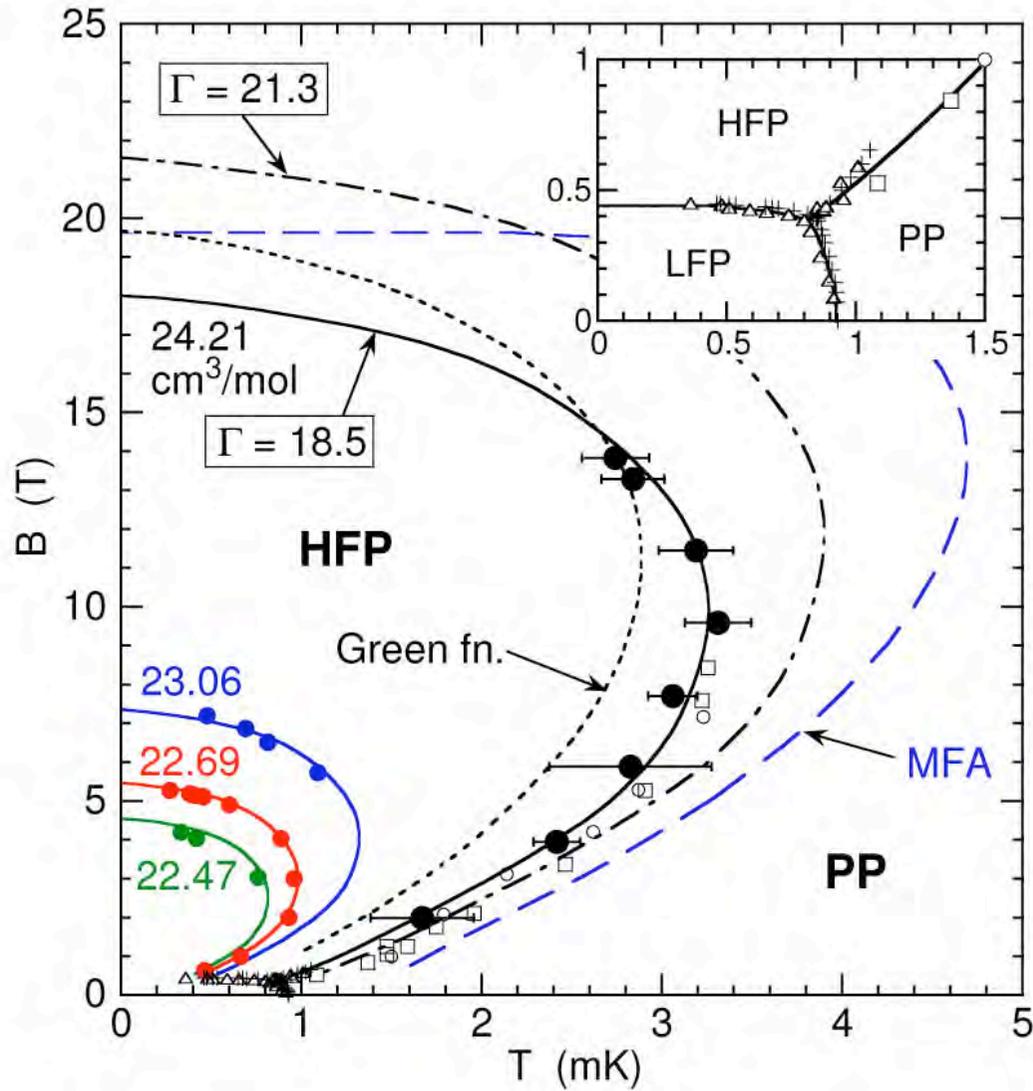
Mizusaki et al. (J. Low Temp. Phys. **127**, 1-10 (2002))

The CNAF Phase of Solid ^3He

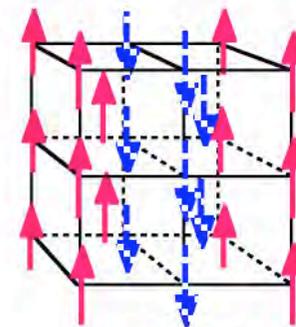


Jack Hetherington

Thanks to Hiroshi Fukuyama

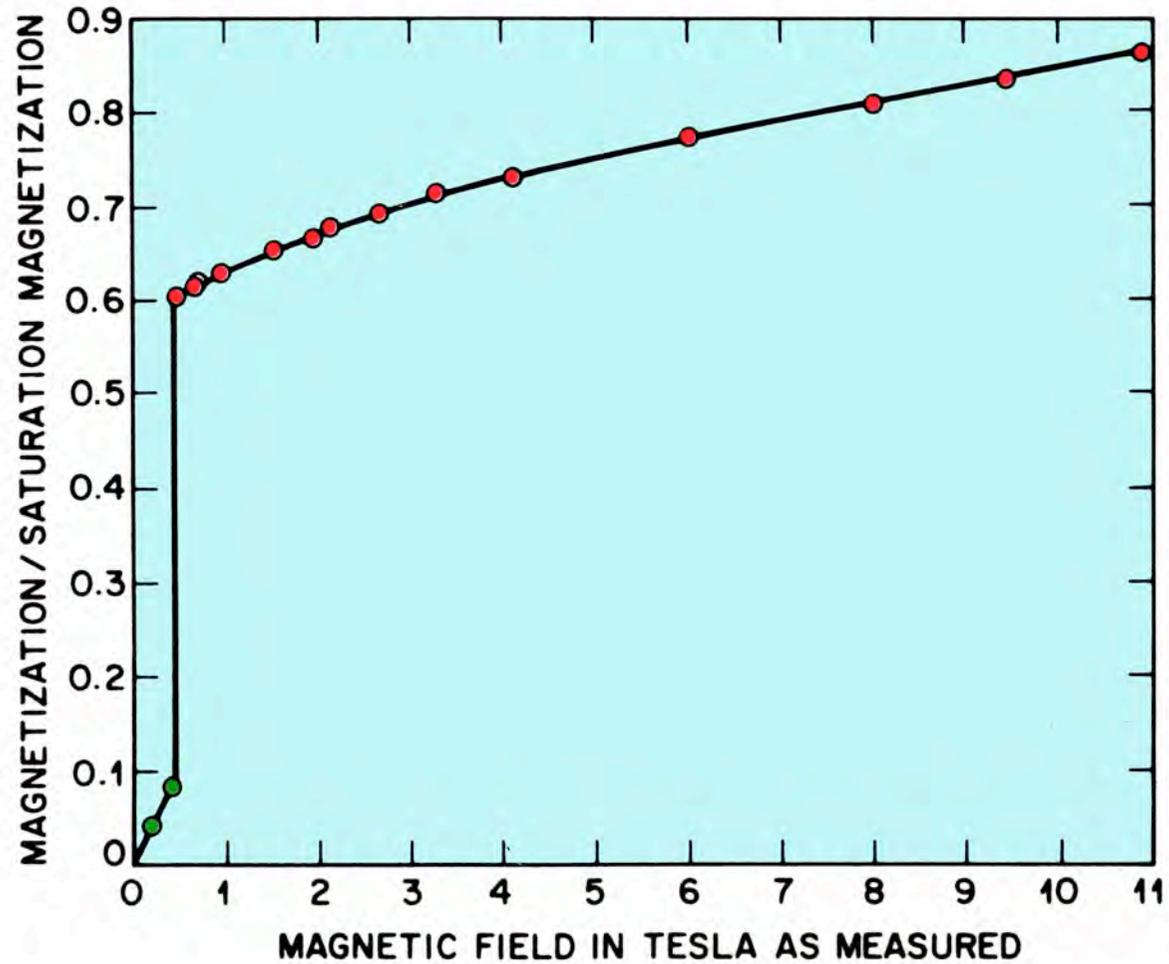


CNAF phase (HFP)

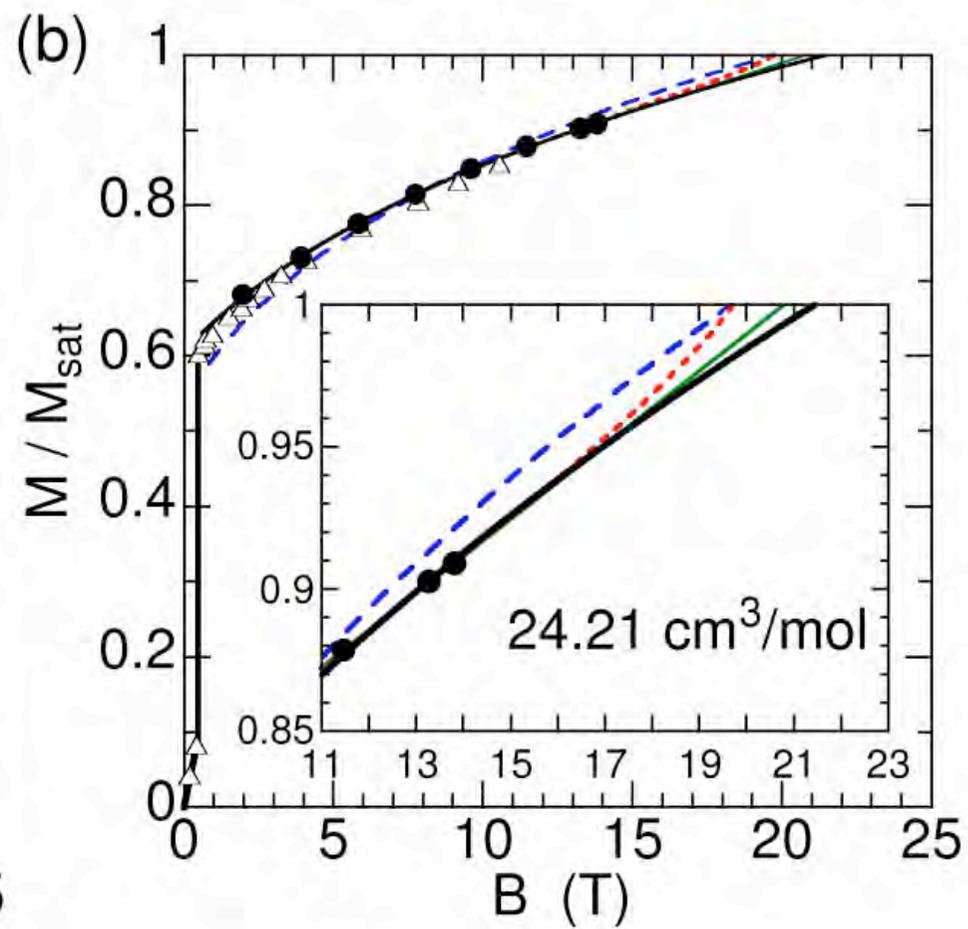
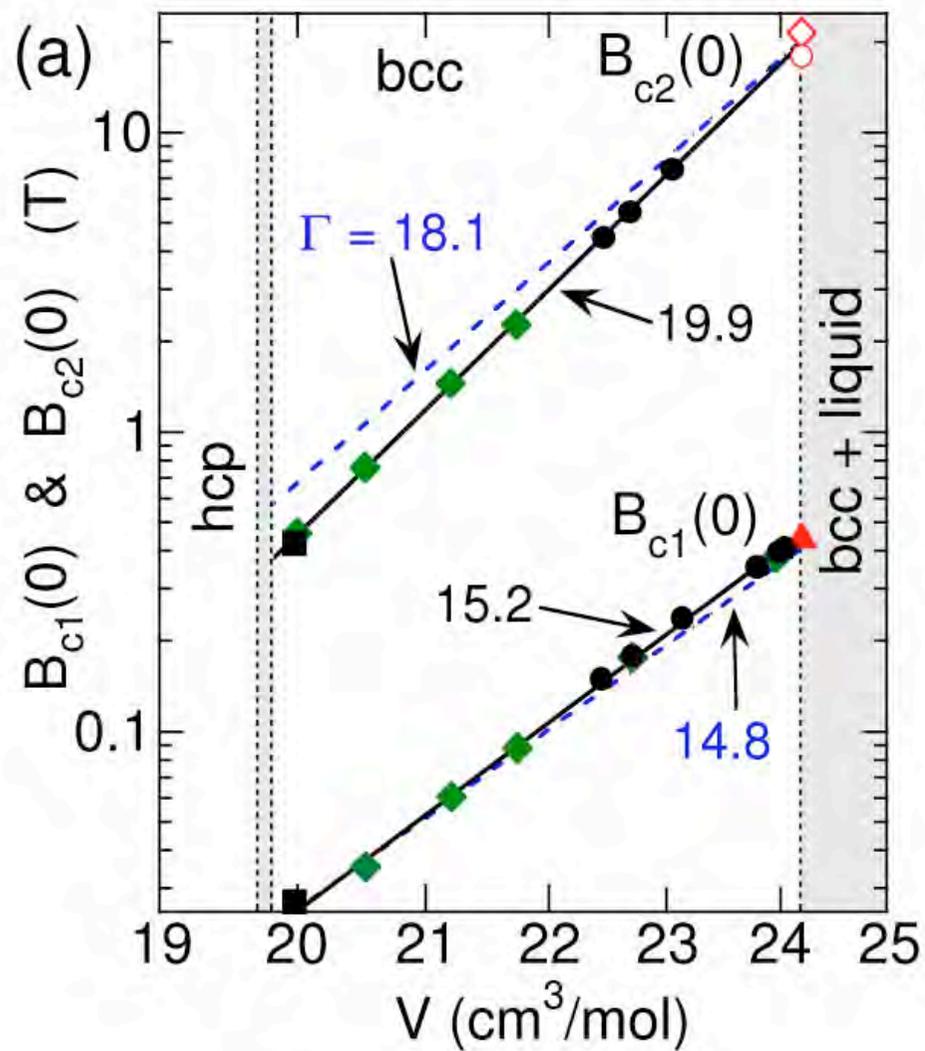


U2D2 phase (LFP)

T=0 Solid ^3He Magnetization



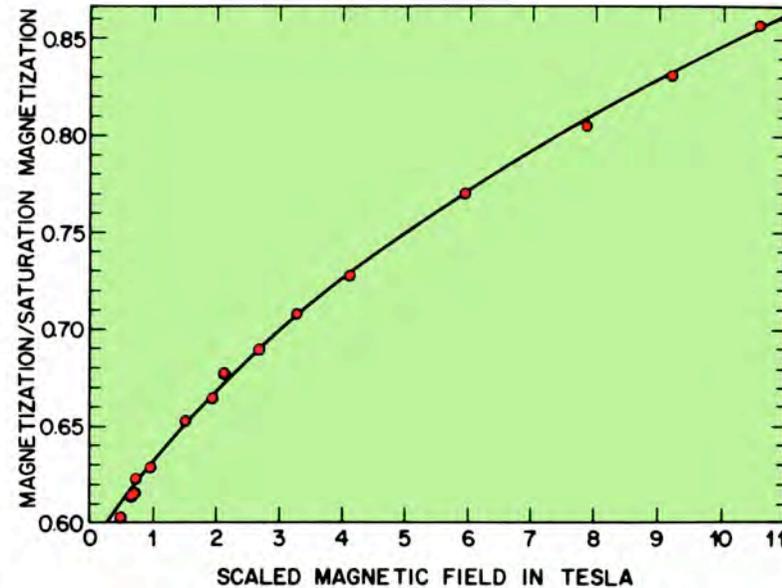
Henri Godfrin and David Ceperley and later Hiroshi Fukuyama



- Change magnetic field by a small well known amount.
- Change in melting pressure just compensates for the change in liquid and solid free energies.
- Calculate solid magnetization
- Express magnetization in terms of magnetic field scaled to change in molar volume from change in P.
- Express magnetization in terms of exchange frequencies.

D.D. Osheroff, H. Godfrin and R. Ruel
Phys. Rev. Letts. 58, 2458 (1987).

CNAF Magnetization Fit



$$H = -(9.1 \pm 0.6)m + (23.8 \pm 2)m^3 + (7.6 \pm 2)m^5$$

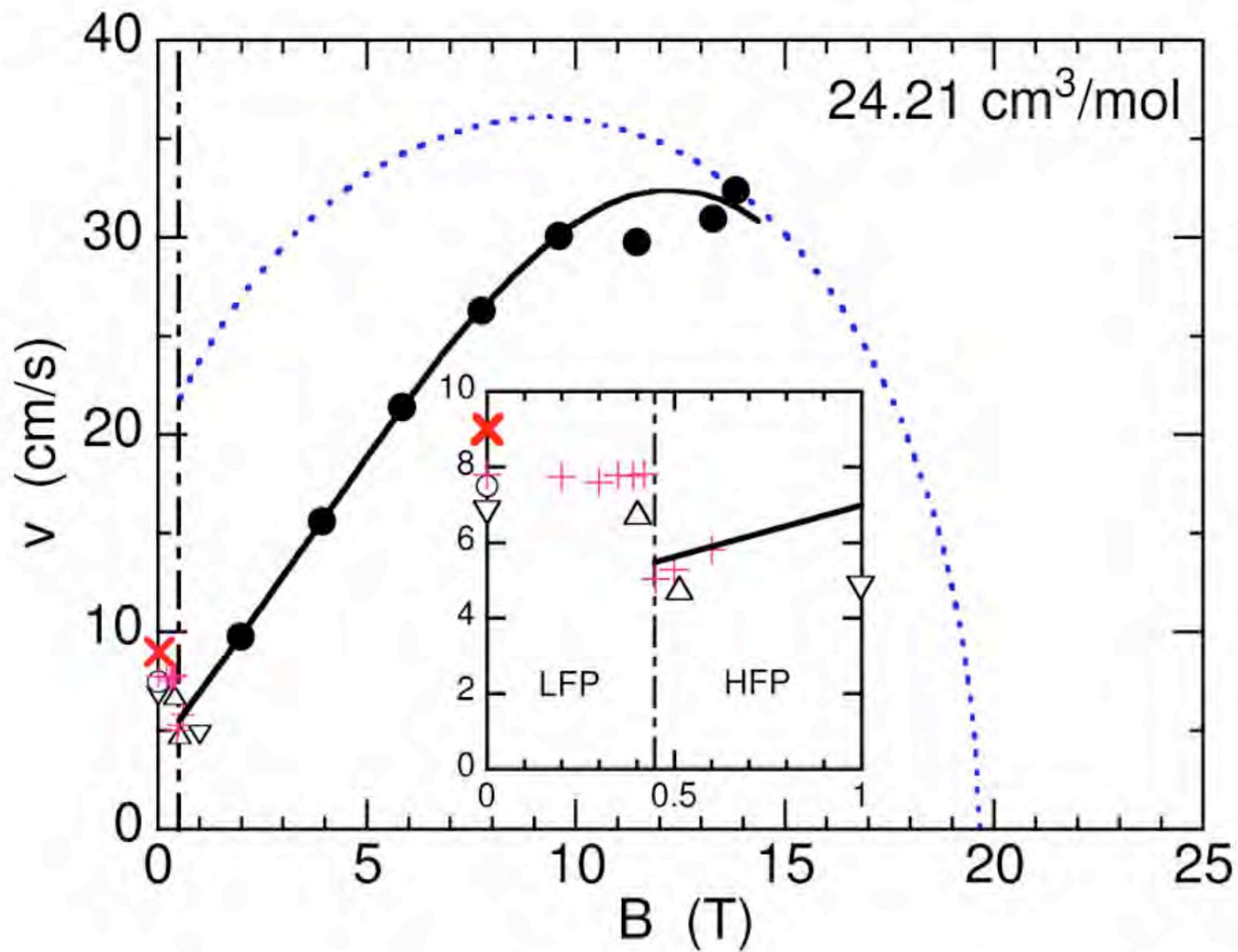
$$H_{C2} = 22.3 \text{ T}$$

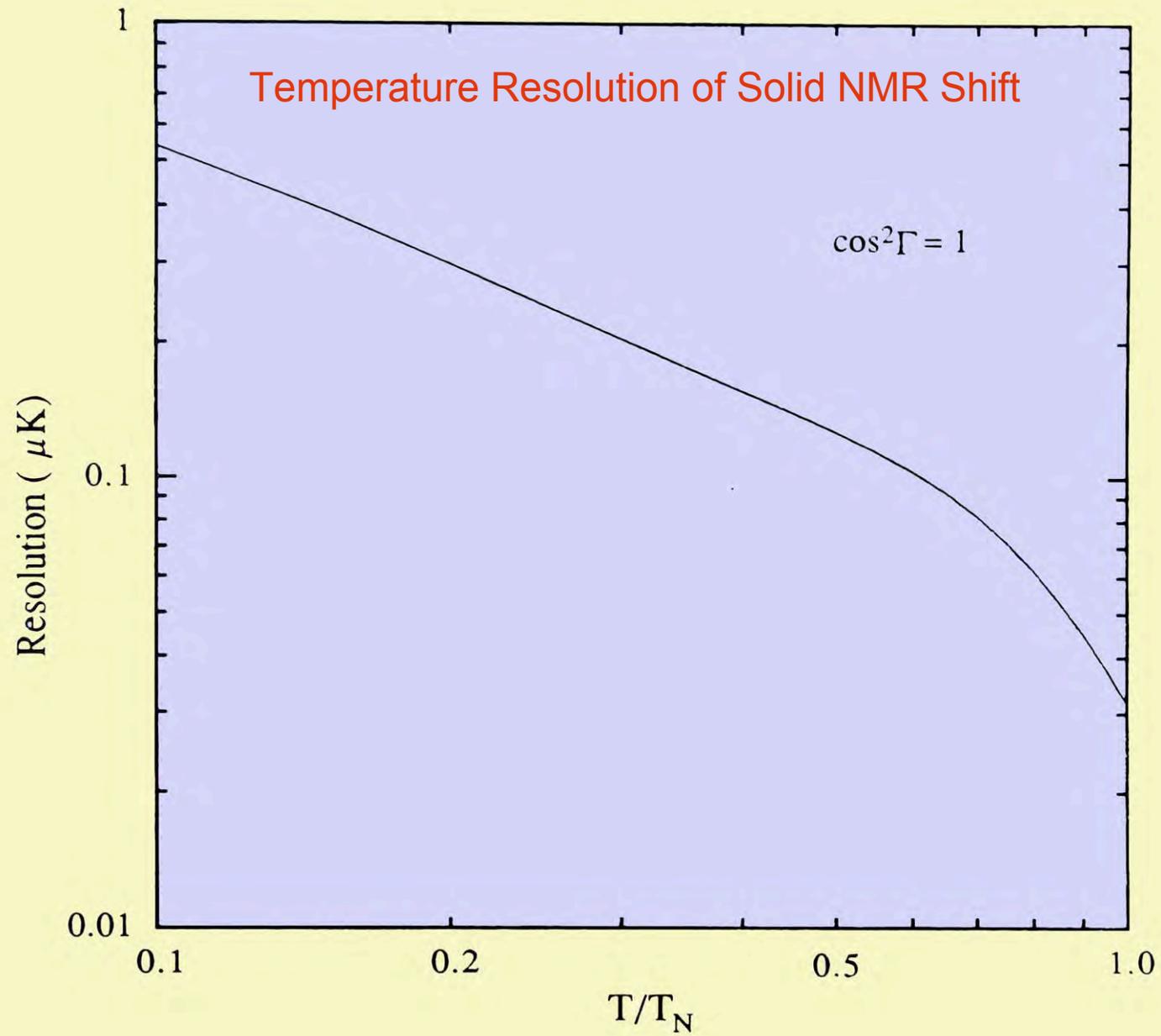
From Ceperley-Jacucci exchange terms:

$$H_{C2} = 20.5 \text{ T}$$

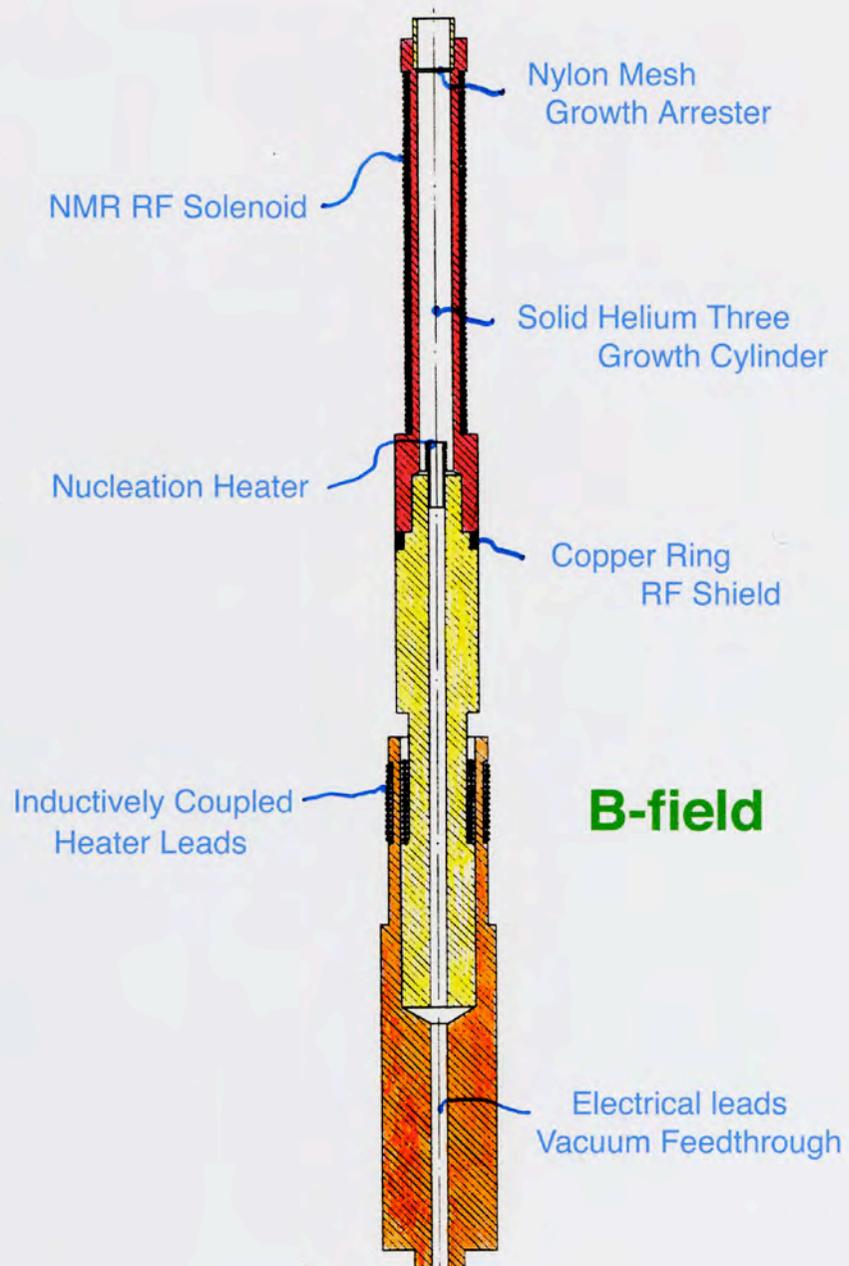
$$H = -(6.0 \pm 1.4)m + (19.7 \pm 2)m^3 + (6.8 \pm 2)m^5$$

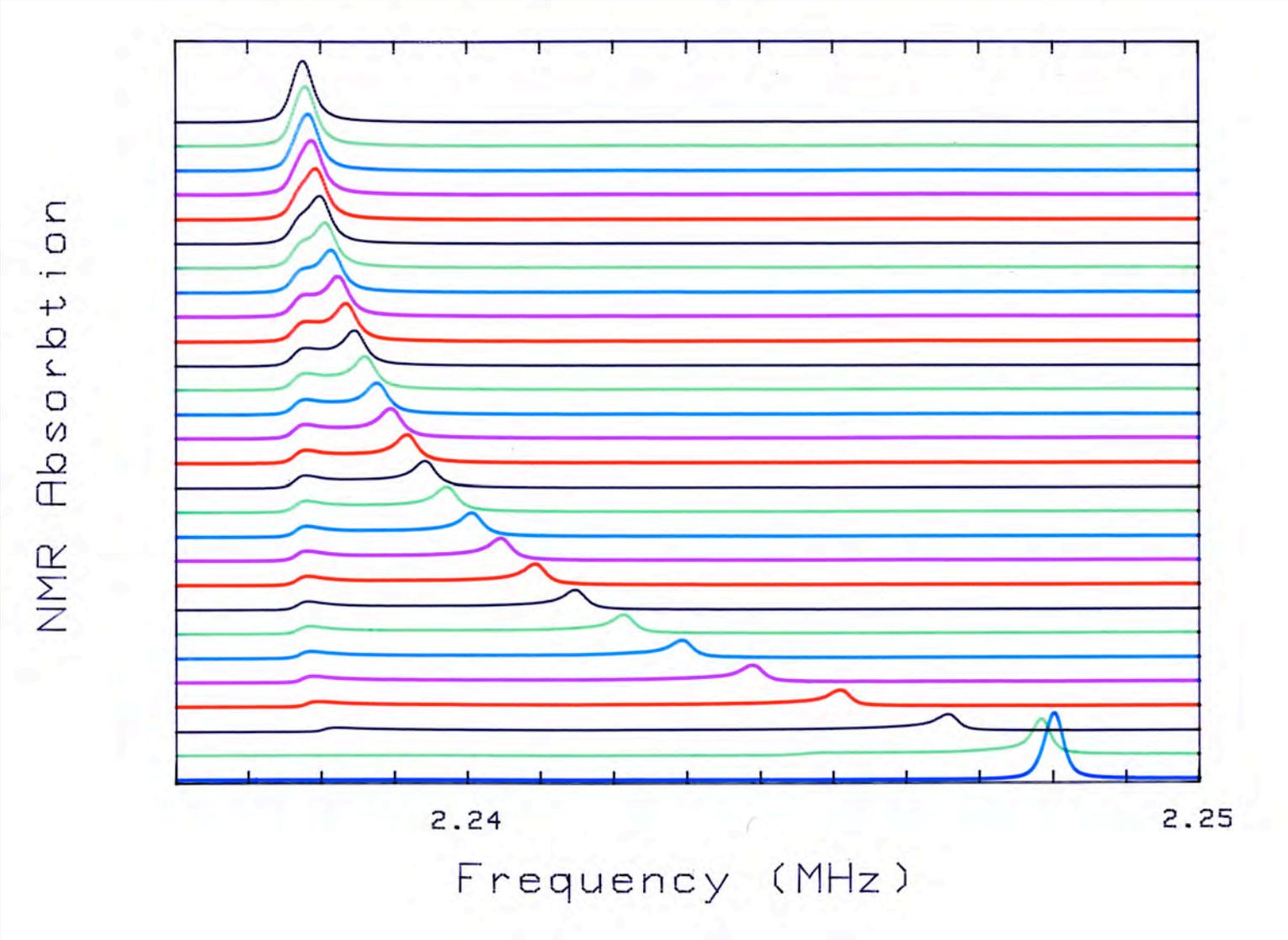
Thanks to Hiroshi Fukuyama

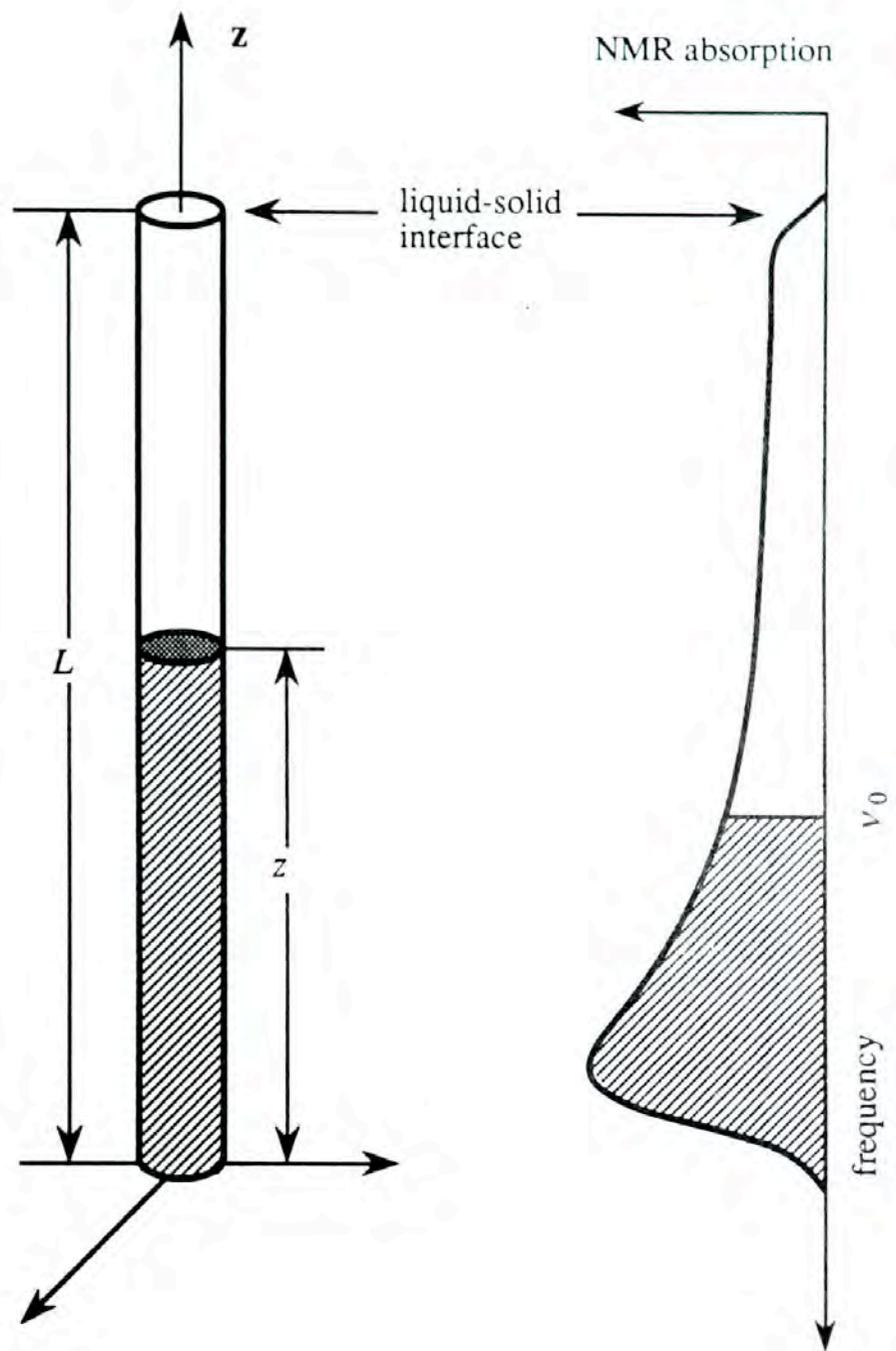


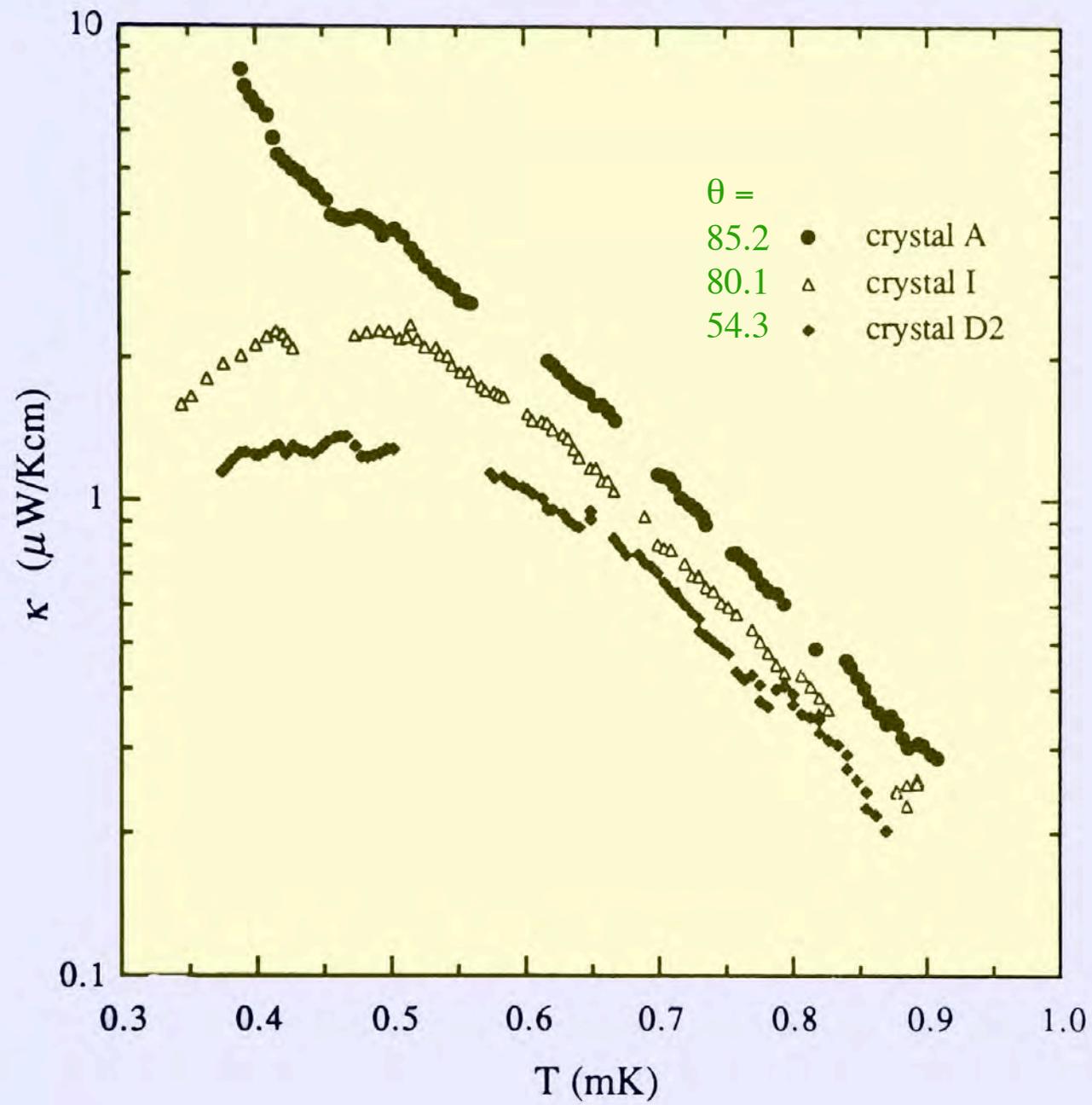


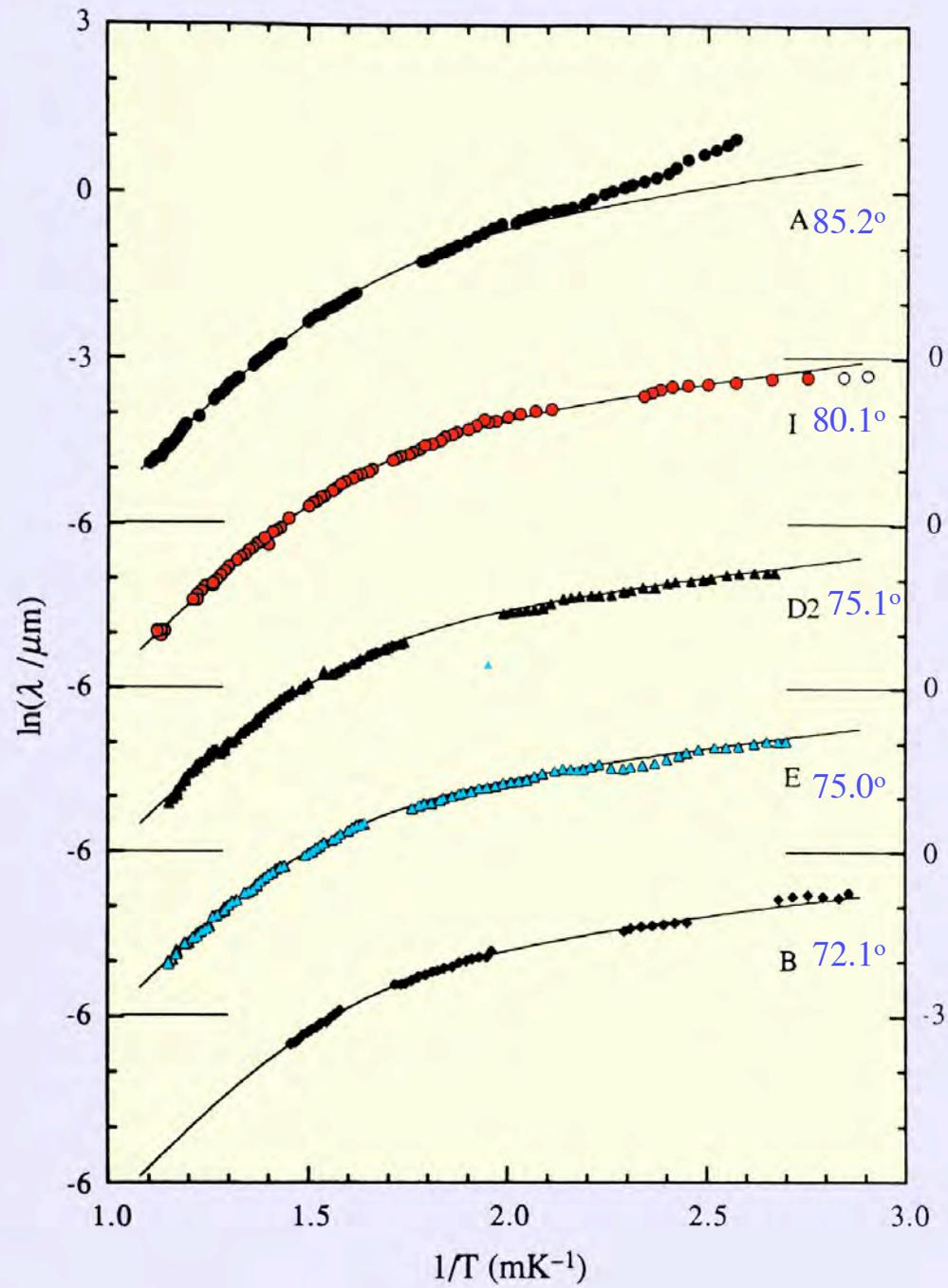
Magnon Diffusivity Cell











Conclusions:

BCC solid ^3He has been a remarkably fruitful system for understanding spin ordering due to particle exchange. Two very different nuclear spin ordered phases coexist over a narrow range of magnetic field. One phase exhibits rich antiferromagnetic resonance spectra while the other allows us to rather directly determine the importance of various ring exchange



*Thank You
for your
attention!*