# Spin-charge interplay in frustrated itinerant systems

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Chirality-driven mass enhancement in the kagome Hubbard model (M. Udagawa and YM, Phys. Rev. Lett. **104**, 106409 (2010); *Fig.1 in Abstract*)

- Phase competition in the pyrochlore double-exchange model (YM and N. Furukawa, Phys. Rev. Lett. **104**, 106407 (2010); *Fig.2 in Abstract*)
- Partial Kondo screening in frustrated Kondo models [**P15** K. Nakamikawa] (K. Nakamikawa, Y. Yamaji, M. Udagawa, and YM, in preparation)
- Non-coplanar order and anomalous Hall effect in the triangular-lattice double-exchange model (Y. Akagi and YM, in preparation)
- ... and more !

[**P16** H. Ishizuka; **P17** J. Yoshitake; **P22** M. Udagawa; **P43** T. Misawa]

## Chirality-driven mass enhancement in the kagome Hubbard model

Masafumi Udagawa and Yukitoshi Motome



- Introduction
	- heavy *d* electrons: Kondo or correlation + frustration?
- Model and Method
- **Results** 
	- energy hierarchy among charge, spin, and chirality
	- heavy-mass behavior due to the degeneracy associated with chirality
- **Summary**

#### Heavy mass: Conventional Kondo physics

■ heavy-fermion *f*-electron systems: hybrid of conduction electrons and localized moments

large mass renormalization due to screening of localized spins by conduction electrons at Kondo temperature  $T \sim T_K$ 

- $\rightarrow$  release of the spin entropy below  $T_K$
- ➡ specific-heat coefficient:  $\gamma \sim \log 2/T_{\rm K}$

 $\checkmark$  localized moments = entropy reservoir



F. Schlottmann, 1989

### Heavy mass in transition metal oxides: Unconventional mass enhancement?

**□** several examples of heavy *d* electrons LiV<sub>2</sub>O<sub>4</sub>,  $(Y_1,X_0X_0)$ Mn<sub>2</sub>,  $\beta$ -Mn, ...

**u** typical: spinel oxide  $Liv_2O_4$ 

• frustrated pyrochlore lattice of V

• no clear sign of phase transition

characteristic temperature *T*\*~20-30K: heavy mass behavior at lower *T*

controversial on the mechanism of heavy mass behavior: no obvious entropy reservoir

*Kondo ?* V. I. Anisimov *et al*., 1999

*electron correlation + frustration ?*

V. Eyert *et al*., 1999; H. Tsunetsugu, 2002; Y. Yamashita and K. Ueda, 2003, etc.



#### Electron correlation + Frustration: A "folklore"

■ Mott criticality (Brinkman-Rice, Gutzwiller, dynamical mean-field, ...)



√ critical mass enhancement in the paramagnetic solution metallic phase, is due to their inability to the inability to the inability to capture their inability to capture the induced state of the  $\blacksquare$  are qualitatively valid in the Mott insulation in the Mott insulation in the Mott insulation insulating insulation in the Mott insulation in the Mott insulation in the Mott insulation in the Mott insulation in the  $\sim$  potopoooooo $\pm$ 

- $\checkmark$  local spin fluctuation under strong correlation = entropy reservoir with a gap  $\mathbb{R}$  is greater than  $\mathbb{R}$  $\mathbf{r}$ FIG. 33. Phase diagram of the fully frustrated model at half-**Filling. It is possible to move continuous**
- √ Usually, all of these are masked by the symmetry breaking  $\mathcal{D}$ y the opservation that in the  $\mathcal{D}$ consistency relation, Eq. (221) implies that '(#+i0<sup>+</sup>  $\Omega$  inside the gap, except for a set of (-function <sup>p</sup>iece in Im' at #=0, with the other since at high temperature the transition becomes a **v** USUAlly, all OI these are rilas

"Folklore": Frustration suppresses the symmetry breaking and rejuvenates the mass enhancement hidden in the 'bare' paramagnetic state. etry breaking and rejuvenates na inaraj naramann the impurity model, we have to deal with an impurity  $\mathsf{H}\cap\mathsf{C}\mathsf{H}\cap\mathsf{H}\cap\mathsf{H}$  $\overline{\phantom{a}}$   $\overline{\phantom{a}}$   $\overline{\phantom{a}}$ approximate location of the actual first-order transition line.  $\Box$  "Folklore": Frustration suppres second-order points.

#### Caveat... id.<br>G

Even in the paramagnetic solution, the quasi-particle peak is fragile against spatial correlations. results for the double occupancy *D*occ % h*ni*"*ni*#i at various **the Even in the paramagnetic solutic** creases as *U* increases, indicating the development of local as the quasi-particle peak is frague

lattice Hubbard model at half filling. In Fig. 2, we show the





On the other hand, in the frustrated case... dependence of the double of the double of the double of the double occupancy in  $\mathcal{L}$  $\Box$  On the other hand,  $\frac{1}{2}$  in the structure is the square lattice case, where the square is  $\frac{1}{2}$  $\Box$  and  $\Box$  in the must also developed and a pseudo-



What is the 'true' role of frustration?

#### **Objectives**

to clarify the role of frustration in correlated metals

- ๏ secondary role, just to suppress the spatial correlations ?
	- $\blacktriangleright$  If yes, mass enhancement occurs in the energy scale of spin  $\sim$  *J*
- ๏ The answer is NO ! (as we will see later)
	- ➡ mass enhancement occurs at much smaller energy scale
- ๏ What determines the smaller new energy scale ?
	- $\rightarrow$  emergent degree of freedom under frustration  $+$  correlation

to explore the new mechanism of quasi-particle mass enhancement

#### Model and Method

Hubbard model on the kagome lattice at half filling

$$
\mathcal{H} = -t \sum_{\langle ij \rangle \sigma} (c_{i\sigma}^{\dagger} c_{j\sigma} + \text{h.c.}) + U \sum_{i} n_{i\uparrow} n_{i\downarrow}
$$

- ๏ a minimal model including both electron correlation and frustration (Y. Imai *et al*., 2003; N. Bulut *et al*., 2005; T. Ohashi *et al*., 2006; B. H. Bernhard *et al*., 2007)
- ๏ Mott transition at *U*c~8.3*t* (T. Ohashi *et al*., 2006)

cluster extension of the dynamical mean-field theory

- ๏ mapping to cluster impurity models (3 or 9 sites)
- ๏ impurity problem solver: continuous-time quantum Monte Carlo method (E. Gull *et al*., 2008)



#### Result: Heavy-fermion behavior phase close to the critical points of the critical points. Fig.  $\alpha$ quasiparticles persist up to the transition point (*U=W* %

dogap formation, which is consistent with the *U* and *T*



To identify the relevant degree of freedom, we calculate  $\bigcirc$ density matrix = prob. distribution of quantum mechanical states  $\rho_\Psi = \frac{1}{Z} \text{Tr}|\Psi\rangle\langle\Psi|e^{-\beta \mathcal{H}}$ 

#### Result: Spin chirality degree of freedom



#### Result: Crossover temperatures



#### Result: Specific heat and entropy



- charge-spin-chirality separation
	- broad hump in  $C_v$  at  $T \sim T_{charge}$
	- entropy~log8 at  $T \sim T_{spin}$
	- sharp peak in *C*v and entropy~log4 at *T*~*T*chirality
- chirality-driven mass enhancement
	- specific-heat coefficient:

 $\gamma \simeq \frac{1}{3} \log 4/T_{\rm chiral}$ 

#### **Discussion**

"Folklore" scenario Frustration just suppresses magnetic LRO **■ heavy mass due to spin entropy** Present mechanism

*T*spin *T*charge

Frustration brings about an emergent degree of freedom, chirality

**■ heavy mass due to spin chirality** 



heavy-mass behavior:  $\bigcirc$ 

crossover from highly-symmetric local state to renormalized Fermi liquid

emergent composite objects with high local symmetry  $\begin{array}{ccc} & \end{array}$   $\begin{array}{ccc} & \end{array}$  chirality, multipole, etc.



### **Summary**

M. Udagawa and YM, Phys. Rev. Lett. **102**, 106409 (2010)

- cellular DMFT study of correlated metallic region in the kagome Hubbard model ✓ continuous-time QMC
	- ✓ cluster-size dependence
- Emergent degree of freedom, chirality, plays a decisive role at low *T*.
	- ✓ energy hierarchy
	- ✓ sharp peak in the specific heat
	- ✓ mass enhancement
- Our results uncover an intensive role of geometrical frustration in correlated metal (not secondary just to suppress LRO).  $_{0.0}$

